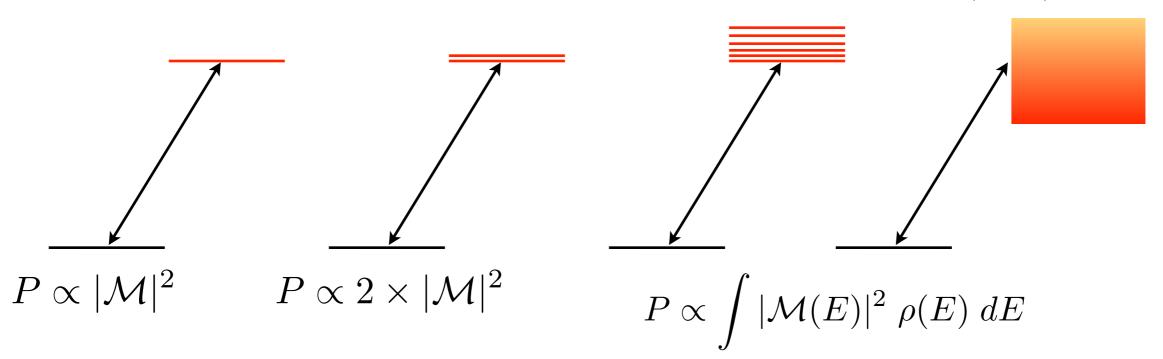
LECTURE 4: PHASE SPACE IN DECAYS

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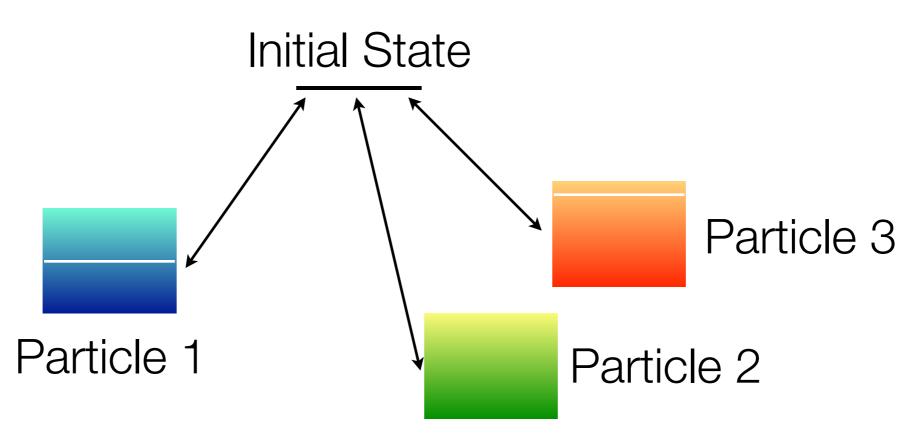
GOLDEN RULE

- Fermi's golden rule states that the probability of a transition in quantum mechanics is given by the product of:
 - The absolute value of the matrix element (a k a amplitude) squared
 - The available density of states. $P \propto |\mathcal{M}|^2 imes
 ho$



- Typically a decay of a particle into states with lighter product masses has more "phase space" and more likely to occur.
- Let's see how to calculate the phase space
 - we'll learn how to calculate amplitudes later.

PRODUCT OF PHASE SPACE

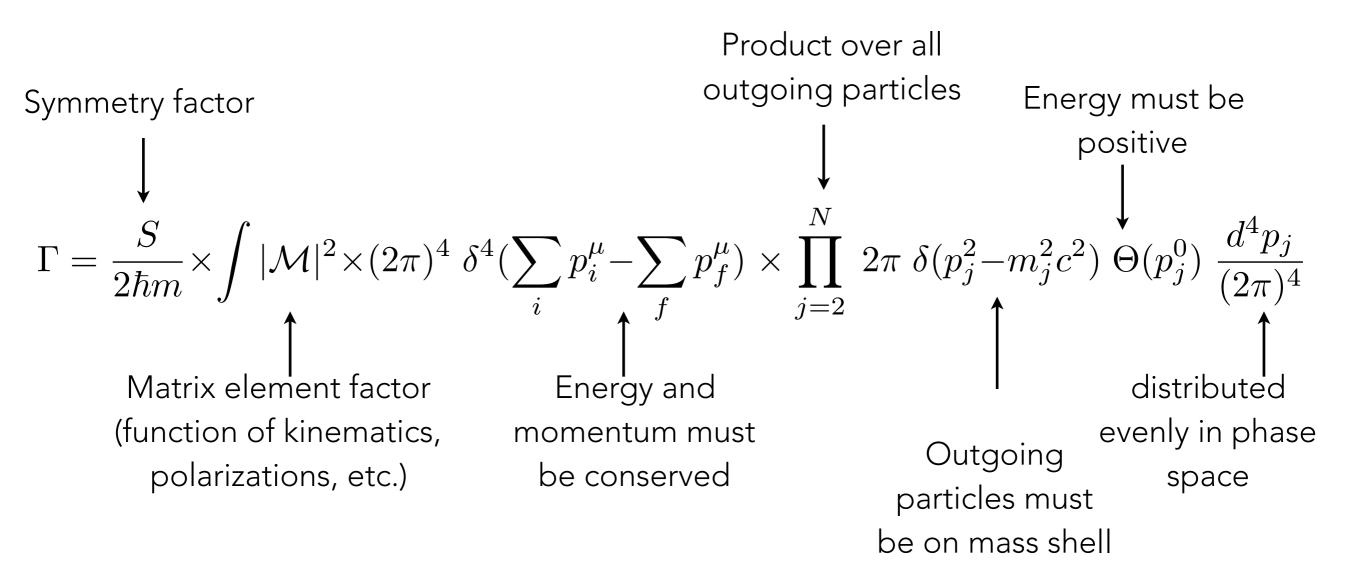


- What is net phase space for the particle 1,2,3 to end up in particular places?
 - 0 if energy and momentum are not conserved
 - 0 if particles are not on "mass shell"
- Otherwise, the product of the individual phase spaces:

$$\rho = \rho_1(p_1^\mu) \times \rho_2(p_2^\mu) \times \rho_3(p_3^\mu)$$

integral extends over region satisfying kinematic constraints $\rho_{tot} = \int_{allowed} \frac{d^4p_1}{(2\pi)^4} \frac{d^4p_2}{(2\pi)^4} \frac{d^4p_3}{(2\pi)^4} \rho_1(p_1^{\mu})\rho_2(p_2^{\mu})\rho_3(p_3^{\mu})$

PHASE SPACE IN DECAYS



- Complicating looking, but represents a basic statement:
 - apart from matrix element, phase space is distributed evenly among all particles subject to mass requirements, E/p conservation
 - "dynamics" like parity violation, etc. incorporated into matrix element.

THE SYMMETRY FACTOR:

• Consider the integration of phase space for two particles in the final state where the particles are of the same species.

$$\int \frac{d^4 p_1}{(2\pi)^4} \frac{d^4 p_2}{(2\pi)^4}$$

- At some point, say, there will be a configuration where $p_1 = K_1$ and $p_2 = K_2$
 - Since the particles are identical, we should also have the reverse case:
 - p₁= K₂, p₂ = K₁
 - the integral will contain both cases separately.
 - However, in quantum mechanics, the identicalness of particles of the same species means that these are the same state and we have double counted.
 - We need to add a factor of 1/2 to the phase space
- Likewise, for n identical particles in the final state, we need a factor of 1/n!

THE GOLDEN RULE: 2-BODY DECAY

$$\Gamma = \frac{S}{2\hbar m_1} \times \int |\mathcal{M}|^2 \times (2\pi)^4 \delta^4(p_1 - p_2 - p_3)$$

$$\times (2\pi) \cdot \delta(p_2^2 - m_2^2 c^2) \Theta(p_2^0) \times (2\pi) \cdot \delta(p_3^2 - m_3^2 c^2) \Theta(p_3^0)$$

$$\frac{d^4 p_2}{(2\pi)^4} \frac{d^4 p_3}{(2\pi)^4}$$

Let's integrate over overall outgoing particle phase space to get the total decay rate $d^4p \equiv dp^0 dp^1 dp^2 dp^3$

• Start with the phase space factors: $d^{-}p^{-1}$

$$\delta(p^2 - m^2 c^2) = \delta((p^0)^2 - \vec{p}^2 - m^2 c^2)$$

$$\delta(p^2 - m^2 c^2) = \frac{1}{2 \times \sqrt{\vec{p}^2 + m^2 c^2}} \left[\delta(p^0 - \sqrt{\vec{p}^2 + m^2 c^2}) + \delta(p^0 + \sqrt{\vec{p}^2 + m^2 c^2}) \right]$$

• Ignore the 2nd $\boldsymbol{\delta}$ function since $\Theta(p^0)$ will be 0 whenever p^0 is negative

$$\delta(p^2 - m^2 c^2) = \frac{1}{2 \times \sqrt{\vec{p}^2 + m^2 c^2}} \delta(p^0 - \sqrt{\vec{p}^2 + m^2 c^2})$$
$$p_0 \Rightarrow \sqrt{\vec{p}^2 + m^2 c^2}$$

ENFORCING ENERGY/MOMENTUM CONSERVATION

• Now integrate over p_{3}^{0} and p_{2}^{0} using the previous relations

$$\Gamma = \frac{S}{2\hbar m_1} \times \int |\mathcal{M}|^2 \times (2\pi)^4 \,\delta^4(p_1^{\mu} - p_2^{\mu} - p_3^{\mu})$$

$$\Gamma = \frac{S}{32\pi^2 \hbar m_1} \times \int |\mathcal{M}|^2 \times \frac{\delta^4 (p_1^{\mu} - p_2^{\mu} - p_3^{\mu})}{\sqrt{\vec{p_2}^2 + m_2^2 c^2} \sqrt{\vec{p_3}^2 + m_3^2 c^2}} d^3 \vec{p_2} d^3 \vec{p_3}$$

DECAY AT REST:

- Decompose the product delta function (particle 1 at rest) $\delta^4(p_1^{\mu} - p_2^{\mu} - p_3^{\mu}) \Rightarrow \delta(m_1c - \sqrt{\vec{p_2}^2 + m_2^2c^2} - \sqrt{\vec{p_3}^2 + m_3^2}) \ \delta^3(-\vec{p_2} - \vec{p_3})$
- Perform the d³p₃ integral

$$\Gamma = \frac{S}{32\pi^2 \hbar m_1} \times \int |\mathcal{M}|^2 \times \frac{\delta^4 (p_1^{\mu} - p_2^{\mu} - p_3^{\mu})}{\sqrt{\vec{p_2}^2 + m_2^2 c^2} \sqrt{\vec{p_3}^2 + m_3^2 c^2}} d^3 \vec{p_2} d^3 \vec{p_3} \sqrt{\vec{p_3}^2 + m_3^2 c^2}$$

$$\Gamma = \frac{S}{32\pi^2\hbar m_1} \times \int |\mathcal{M}|^2 \times \frac{\delta(m_1c - \sqrt{\vec{p}_2^2 + m_2^2c^2} - \sqrt{\vec{p}_2^2 + m_3^2c^2})}{\sqrt{\vec{p}_2^2 + m_2^2c^2}\sqrt{\vec{p}_2^2 + m_3^2c^2}} d^3\vec{p}_2$$

INTEGRAL IN SPHERICAL COORDINATES

$$d^{3}\vec{p_{2}} \Rightarrow d\phi \ d\cos\theta \ |p_{2}|^{2}dp_{2}$$
Assume no dependence of M on p
$$\Gamma = \frac{S}{32\pi^{2}hm_{1}} \times \int d\phi \ d\cos\theta \ |\mathbf{p}_{2}|^{2} \ d|\mathbf{p}_{2}||\mathcal{M}|^{2} \times \frac{\delta(m_{1}c - \sqrt{\mathbf{p}_{2}^{2} + m_{2}^{2}c^{2}} - \sqrt{\mathbf{p}_{2}^{2} + m_{3}^{2}c^{2}}}{\sqrt{\mathbf{p}_{2}^{2} + m_{2}^{2}c^{2}} \sqrt{\mathbf{p}_{2}^{2} + m_{3}^{2}c^{2}}}$$

$$\int_{0}^{2\pi} d\phi \rightarrow 2\pi \qquad \int_{-1}^{+1} d\cos\theta \rightarrow 2$$

$$u = \sqrt{\mathbf{p}_{2}^{2} + m_{2}^{2}c^{2}} + \sqrt{\mathbf{p}_{2}^{2} + m_{3}^{2}c^{2}}$$

$$du = \frac{u|\mathbf{p}_{2}|}{\sqrt{\mathbf{p}_{2}^{2} + m_{2}^{2}c^{2}}\sqrt{\mathbf{p}_{2}^{2} + m_{3}^{2}c^{2}}} d|\mathbf{p}_{2}|$$

$$\Gamma = \frac{S}{8\pi\hbar m_{1}} \times \int_{p_{2}=0, u=m_{2}+m_{3}}^{\infty} du \ |\mathcal{M}|^{2} \times \ \delta(m_{1}c - u) \frac{|\mathbf{p}_{2}|}{u}$$
The final integral over u sends u=m_{1}c and makes p_{2} consistent with E conservation

FINAL RESULT: TOTAL TWO-BODY DECAY RATE:

$$\Gamma = \frac{S}{8\pi\hbar m_1^2 c} \times |\mathcal{M}|^2 \times |\vec{p_2}|$$

- We now need to be able to calculate the matrix element ${\cal M}$
 - We'll use Feynman diagrams and the associated calculus to calculate amplitudes for various elementary processes

SUMMARY

• Please read 4.1-4.5 for next time