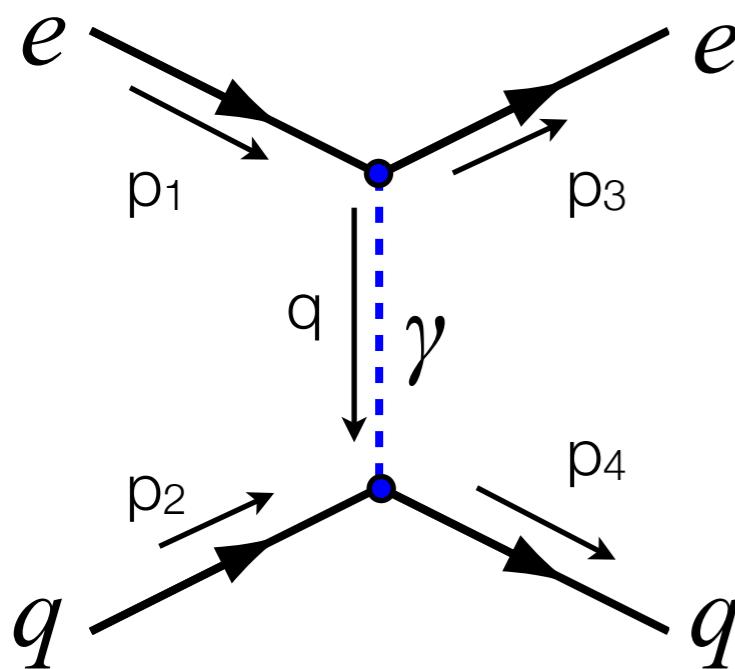




PHYSICS 489/1489

# LECTURE 11: ELECTRON-PROTON SCATTERING

# ELECTRON-QUARK SCATTERING



$$\mathcal{M} = \frac{Qe^2}{(p_1 - p_3)^2} [\bar{u}_3 \gamma^\mu u_1] [\bar{u}_4 \gamma_\mu u_2]$$

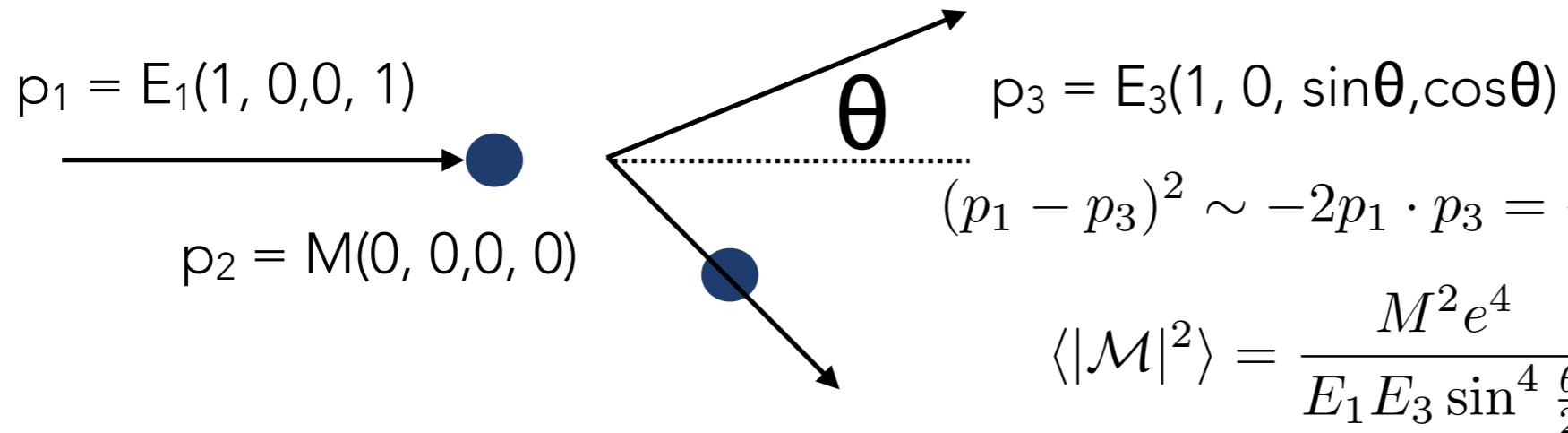
$$|\mathcal{M}|^2 = \frac{Q^2 e^4}{(p_1 - p_3)^4} [\bar{u}_3 \gamma^\mu u_1] [\bar{u}_3 \gamma^\nu u_1]^* [\bar{u}_4 \gamma_\mu u_2] [\bar{u}_4 \gamma_\nu u_2]^*$$

$$\sum_{s_a, s_b} [\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \bar{\Gamma}_2 u_b]^* \\ \Rightarrow \text{Tr} [\Gamma_1 (\not{p}_b + m_b c) \bar{\Gamma}_2 (\not{p}_a + m_a c)]$$

$$\sum_{\text{spins}} |\mathcal{M}|^2 = \frac{e^4}{(p_1 - p_3)^4} \text{Tr} [\gamma^\mu (\not{p}_1 + mc) \gamma^\nu (\not{p}_3 + mc)] \text{Tr} [\gamma_\mu (\not{p}_2 + Mc) \gamma_\nu (\not{p}_4 + Mc)]$$

$$\langle |\mathcal{M}|^2 \rangle = \frac{8Q^2 e^4}{(p_1 - p_3)^4} [(p_1 \cdot p_2)(p_3 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) - (M^2 c^2)(p_1 \cdot p_3)]$$

with  $m \sim 0$



$$(p_1 - p_3)^2 \sim -2p_1 \cdot p_3 = -4E_1 E_3 \sin^2 \frac{\theta}{2} = q^2 \equiv -Q^2$$

$$\langle |\mathcal{M}|^2 \rangle = \frac{M^2 e^4}{E_1 E_3 \sin^4 \frac{\theta}{2}} \left[ \cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right]$$

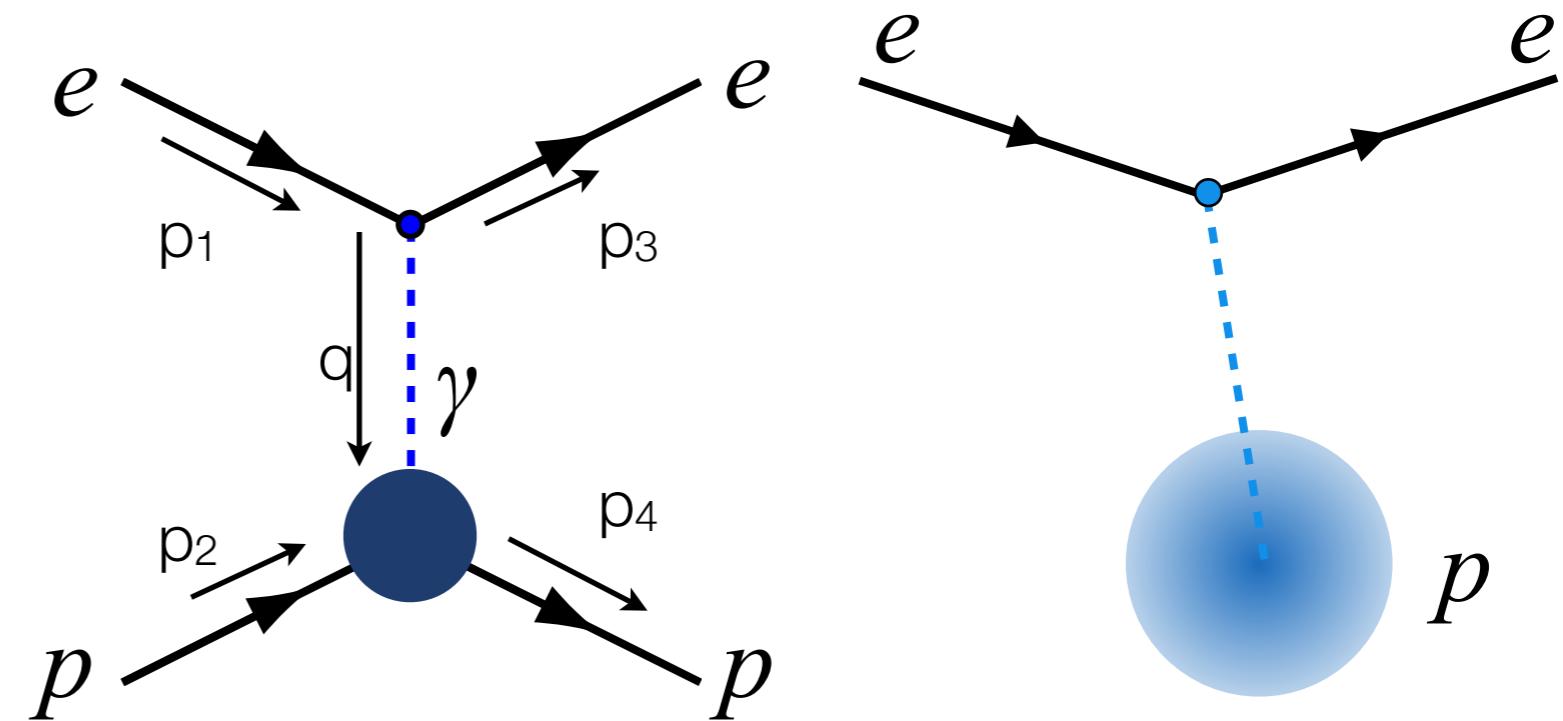
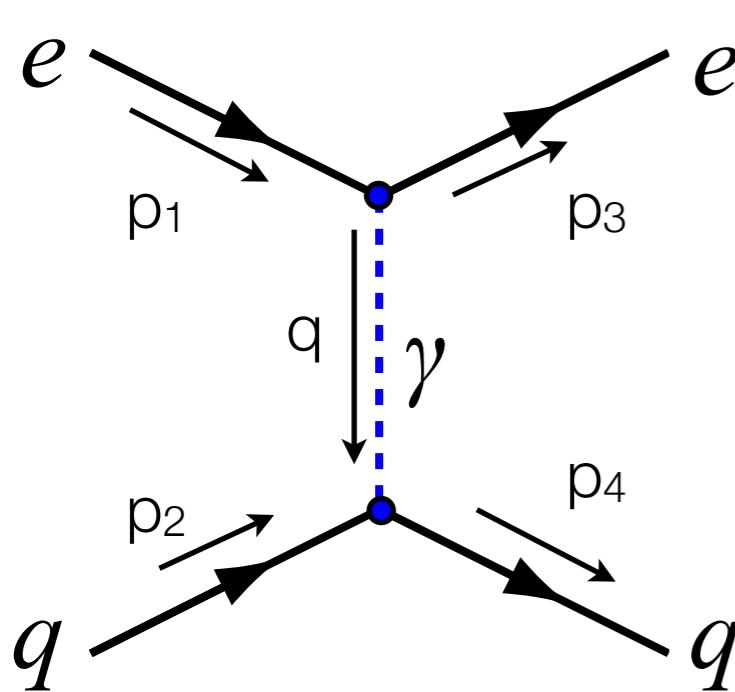
# A FEW NOTES

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left( \frac{E_3}{ME_1} \right)^2 |\mathcal{M}|^2$$
$$= \frac{e^4}{64\pi^2} \frac{1}{E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[ \cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right] = \frac{\alpha^2}{4} \frac{1}{E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[ \cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right]$$

- Note that there is only one degree of freedom in the scattering apart from the incident electron energy  $E_1$ :
  - $E_3 = \frac{E_1 M}{M + E_1(1 - \cos \theta)}$
  - $Q^2 = \frac{2ME_1^2(1 - \cos \theta)}{M + E_1(1 - \cos \theta)}$
  - i.e. if we know one of  $E_3$ ,  $\theta$ , or  $Q^2$ , the others are determined.

# POINT-LIKE PARTICLES

- So far, we have focussed on point particles
  - as far as we know, electrons and quarks do not have any spatial extent
- What if the particle we are dealing with is not fundamental?



# FORM FACTORS: APPROACH 1

- Accounting for spatial distribution of an arbitrary charge distribution:

$$\psi_3(\mathbf{r}) = e^{i(\mathbf{p}_3 \cdot \mathbf{x} - E_3 t)}$$

$$\psi_1(\mathbf{r}) = e^{i(\mathbf{p}_1 \cdot \mathbf{x} - E_1 t)}$$

$$\mathcal{M} = \int d^3\mathbf{r} \psi_3^*(\mathbf{r}) V(\mathbf{r}) \psi_1(\mathbf{r})$$

$$V(\mathbf{r}) = Q \int d^3\mathbf{r}' \frac{\rho(\mathbf{r}')}{4\pi|\mathbf{r} - \mathbf{r}'|}$$

$\mathbf{q} = \mathbf{p}_1 - \mathbf{p}_3$

The diagram illustrates the vectors involved in the calculation. Three momentum vectors,  $\mathbf{p}_1$ ,  $\mathbf{p}_3$ , and  $\mathbf{q} = \mathbf{p}_1 - \mathbf{p}_3$ , are shown originating from a common point. A blue circle represents a charge distribution, with its center at  $\mathbf{r}'$  and radius  $\mathbf{r}$ . The position vector  $\mathbf{R} = \mathbf{r} - \mathbf{r}'$  is also indicated.

$$\mathcal{M} = Q \int d^3\mathbf{r} \int d^3\mathbf{r}' e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{i\mathbf{q} \cdot (\mathbf{r}')} \frac{\rho(\mathbf{r}')}{4\pi|\mathbf{r} - \mathbf{r}'|}$$

$$= Q \int d^3\mathbf{R} \frac{e^{i\mathbf{q} \cdot \mathbf{R}}}{4\pi|\mathbf{R}|} \int d^3\mathbf{r}' \rho(\mathbf{r}') e^{i\mathbf{q} \cdot (\mathbf{r}')}$$

# ANALYSIS

$$\mathcal{M} = Q \left[ \int d^3\mathbf{R} \frac{e^{i\mathbf{q}\cdot\mathbf{R}}}{4\pi|\mathbf{R}|} \right] \left[ \int d^3\mathbf{r}' \rho(\mathbf{r}') e^{i\mathbf{q}\cdot(\mathbf{r}')} \right]$$

- First term is matrix element for a point charge

$$\begin{aligned} \mathcal{M} &= Q \int d^3\mathbf{r} \int d^3\mathbf{r}' e^{i\mathbf{q}\cdot(\mathbf{r}-\mathbf{r}')} e^{i\mathbf{q}\cdot(\mathbf{r}')} \frac{\rho(\mathbf{r}')}{4\pi|\mathbf{r}-\mathbf{r}'|} \\ &= Q \int d^3\mathbf{r} e^{i\mathbf{q}\cdot\mathbf{r}} \frac{1}{4\pi|\mathbf{r}|} \quad \rho(\mathbf{r}') = \delta(\mathbf{0}) \end{aligned}$$

- Express as a modification of the matrix element of a point charge.

$$\mathcal{M} = \mathcal{M}_0 \times F(\mathbf{q}^2)$$

- also consider the distribution of magnetic moment to get another factor

# BACK TO THE CROSS SECTION

$$\left( \frac{d\sigma}{d\Omega} \right)_0 = \frac{\alpha^2}{4} \frac{1}{E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[ \cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right] \quad \tau = \frac{Q^2}{4M^2}$$

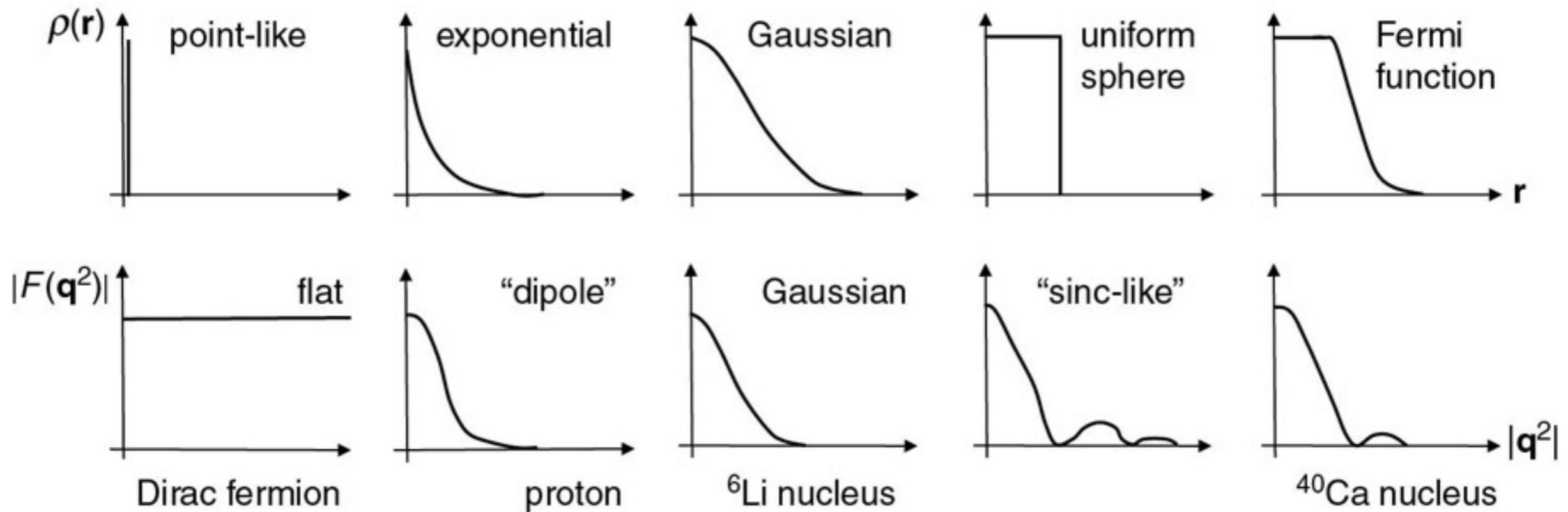
$$\Rightarrow \frac{\alpha^2}{4E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[ \frac{G_E^2(Q^2) + \tau G_M^2(q^2)}{1 + \tau} \cos^2 \frac{\theta}{2} + 2\tau G_M^2(Q^2) \sin^2 \frac{\theta}{2} \right]$$

- Connecting to the previous discussion
- If  $Q^2 \ll M^2$  then to a good approximation  $Q^2 \sim q^2$

$$Q^2 = \mathbf{q}^2 - (E_1 - E_3)^2$$

$$E_1 - E_3 = \frac{Q^2}{2M} \quad \Rightarrow Q^2 \left( 1 + \frac{Q^2}{4M^2} \right) = \mathbf{q}^2$$

# PARMETRIZATIONS



- Measure the form factors as a function of  $q^2$  or  $Q^2$ 
  - effectively measuring the Fourier transform of the charge/magnetic moment of the proton
  - most results use the "dipole" parametrization

$$F(q^2) = \left( \frac{1}{1 + \frac{q^2}{M^2}} \right)^2$$

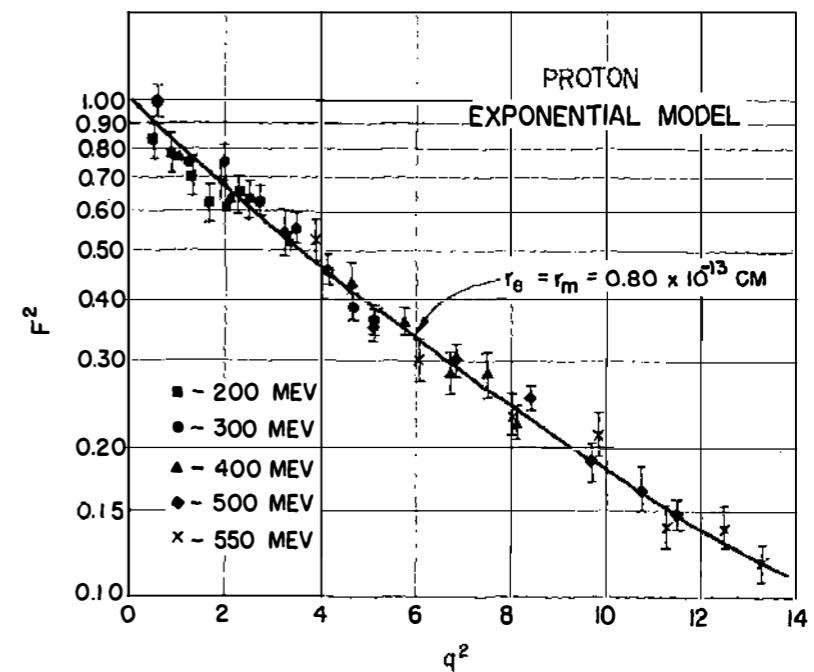


FIG. 8. An example of a model which fits the experimental values of  $F^2$ . This model gives the proton an exponential charge density and an exponential magnetic-moment density with rms radii  $0.80 \times 10^{-13} \text{ cm}$ . Other close models fit equally well.

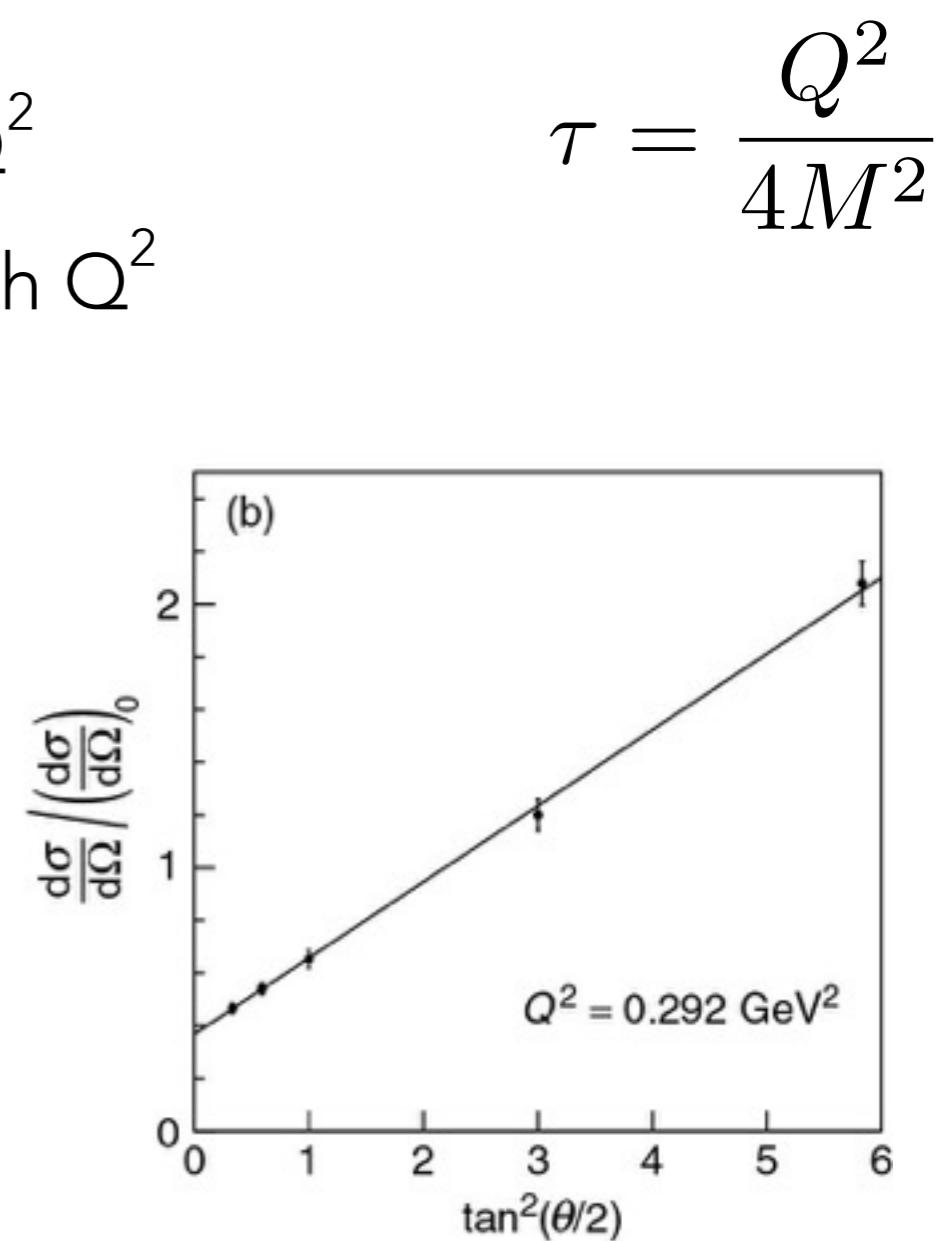
# SEPARATING THE FORM FACTORS

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[ \frac{G_E^2(Q^2) + \tau G_M^2(Q^2)}{1 + \tau} \cos^2 \frac{\theta}{2} + 2\tau G_M^2(Q^2) \sin^2 \frac{\theta}{2} \right]$$

- Electric form factor ( $G_E$ ) is dominant at low  $Q^2$
- Magnetic form factor ( $G_M$ ) is dominant at high  $Q^2$
- General strategy:

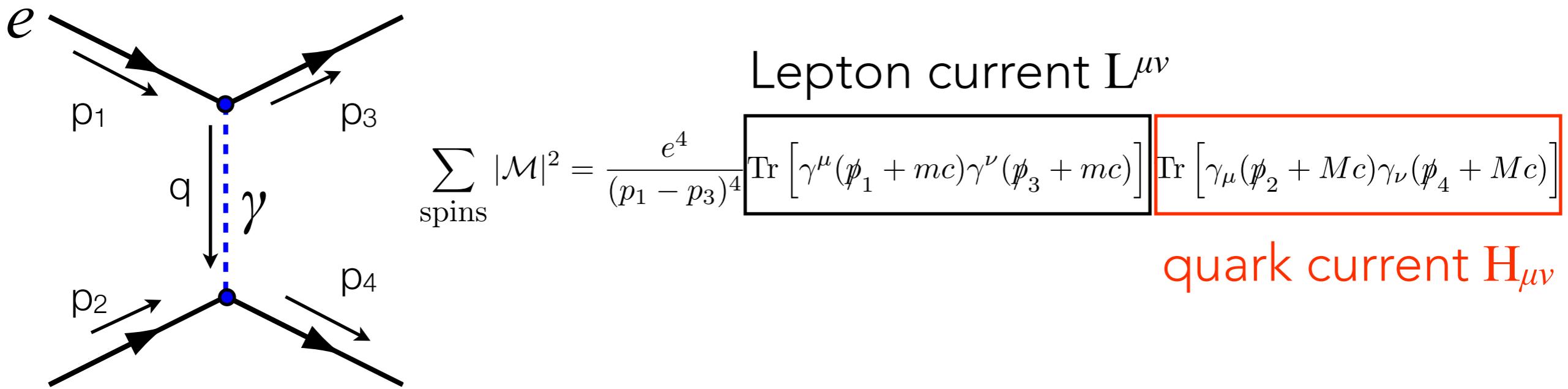
- Vary beam energy
- Measure cross section vs. angle at a fixed  $Q^2$
- compare relative to "Mott" Cross section:

$$\left( \frac{d\sigma}{d\Omega} \right)_0 = \frac{\alpha^2}{4E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \cos^2 \frac{\theta}{2}$$



# FORM FACTORS: APPROACH II

- Another approach to form factors:



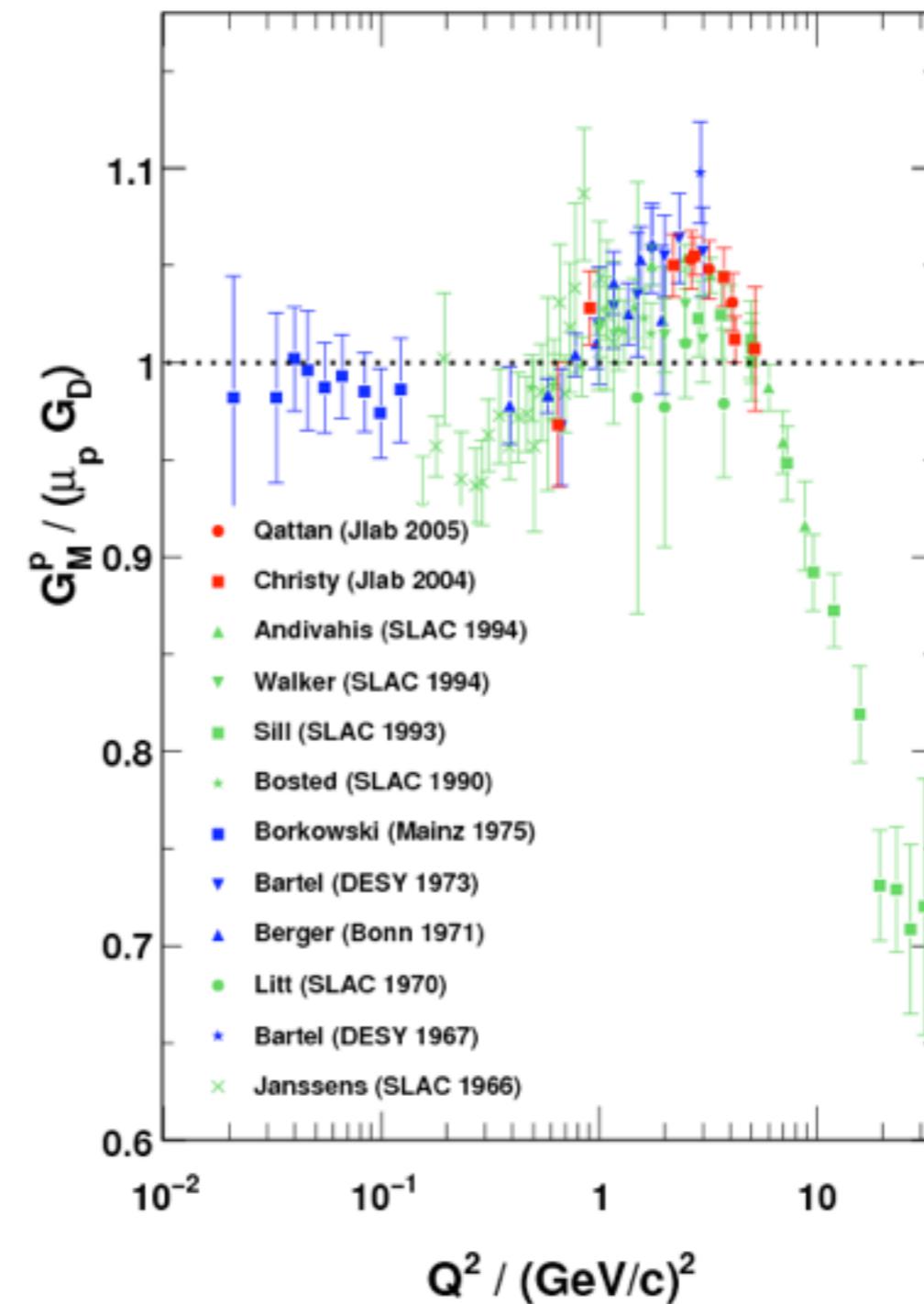
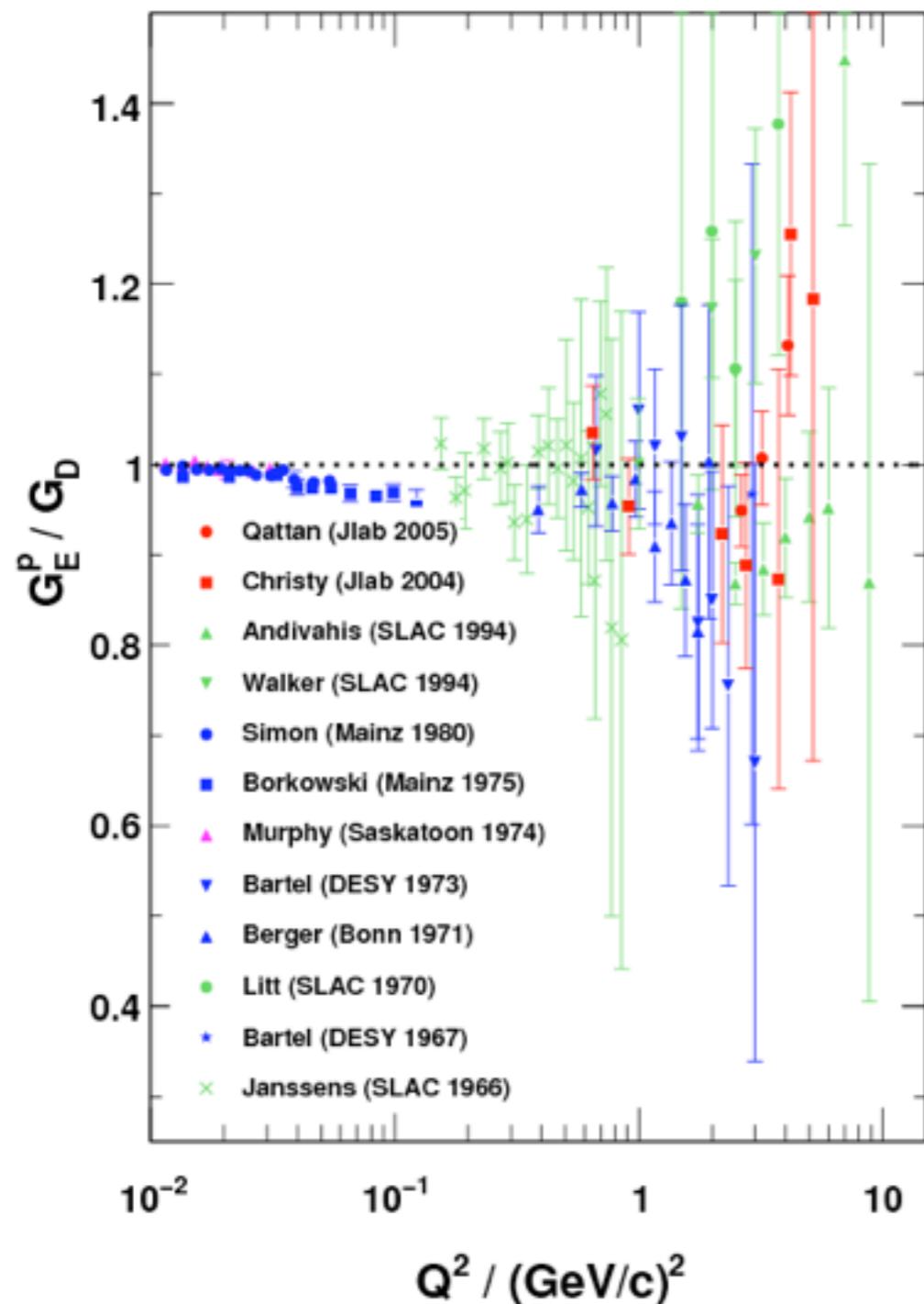
- What form can the hadronic current take?

- it must carry the Lorentz indices  $\mu\nu$
- Lorentz quantities with indices:  $g_{\mu\nu}, p_2, p_4, q$ 
  - $p_4 = p_2 + q$
  - can be a function of Lorentz scale quantities (masses,  $q^2$ )

$$H_{\mu\nu} = H_1(q^2)g_{\mu\nu} + H_2(q^2)p_{2\mu}p_{2\nu} + H_3(q^2)q_\mu q_\nu + H_4(q^2)(p_{2\mu}q_\nu + p_{2\nu}q_\mu)$$

- Other considerations (charge conservation) lead to relations that reduce the number of independent functions to two

# CONTEMPORARY TOPICS



Deviations for dipole for.

Can we obtain a more fundamental understanding of the proton form

# "PARTONS"

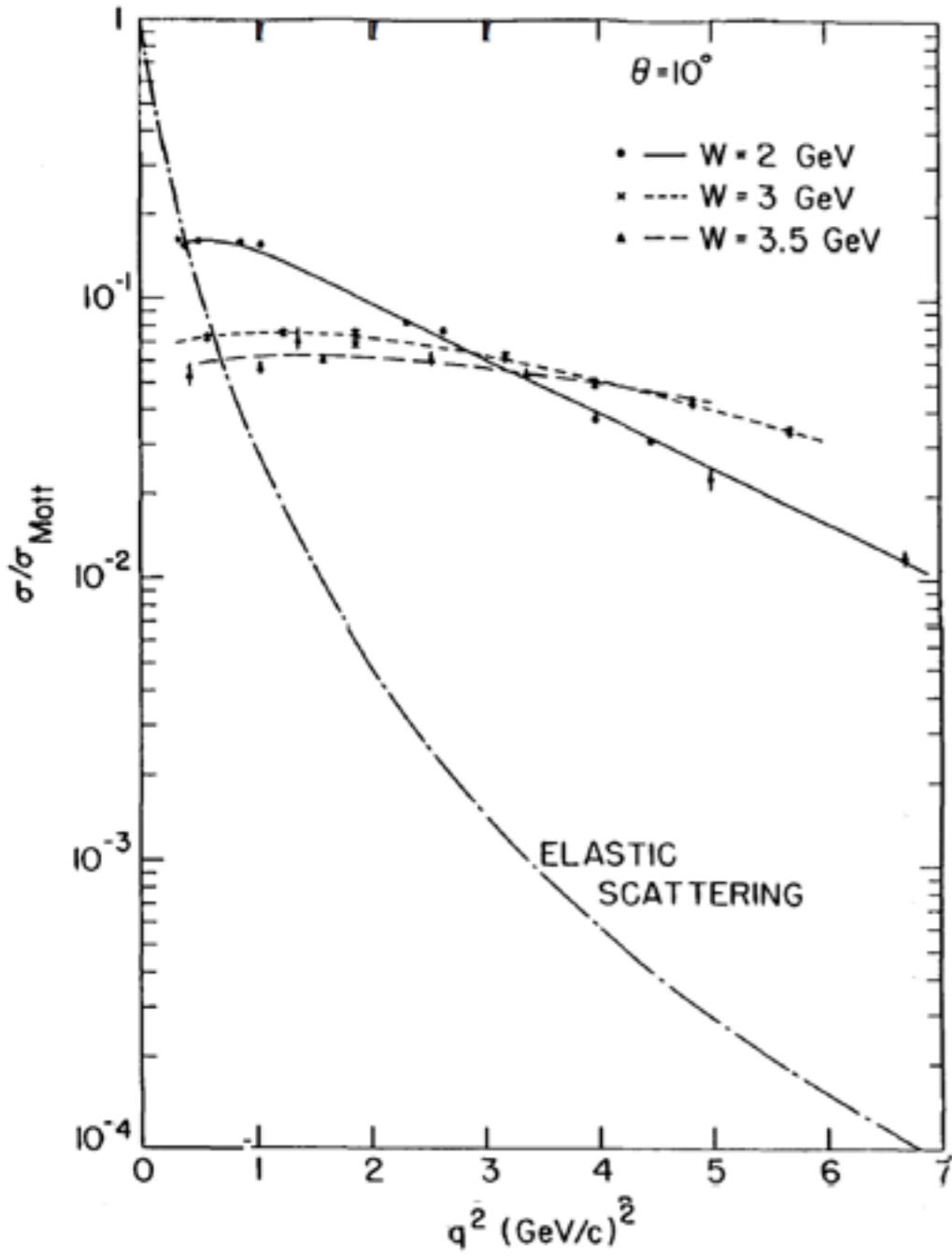
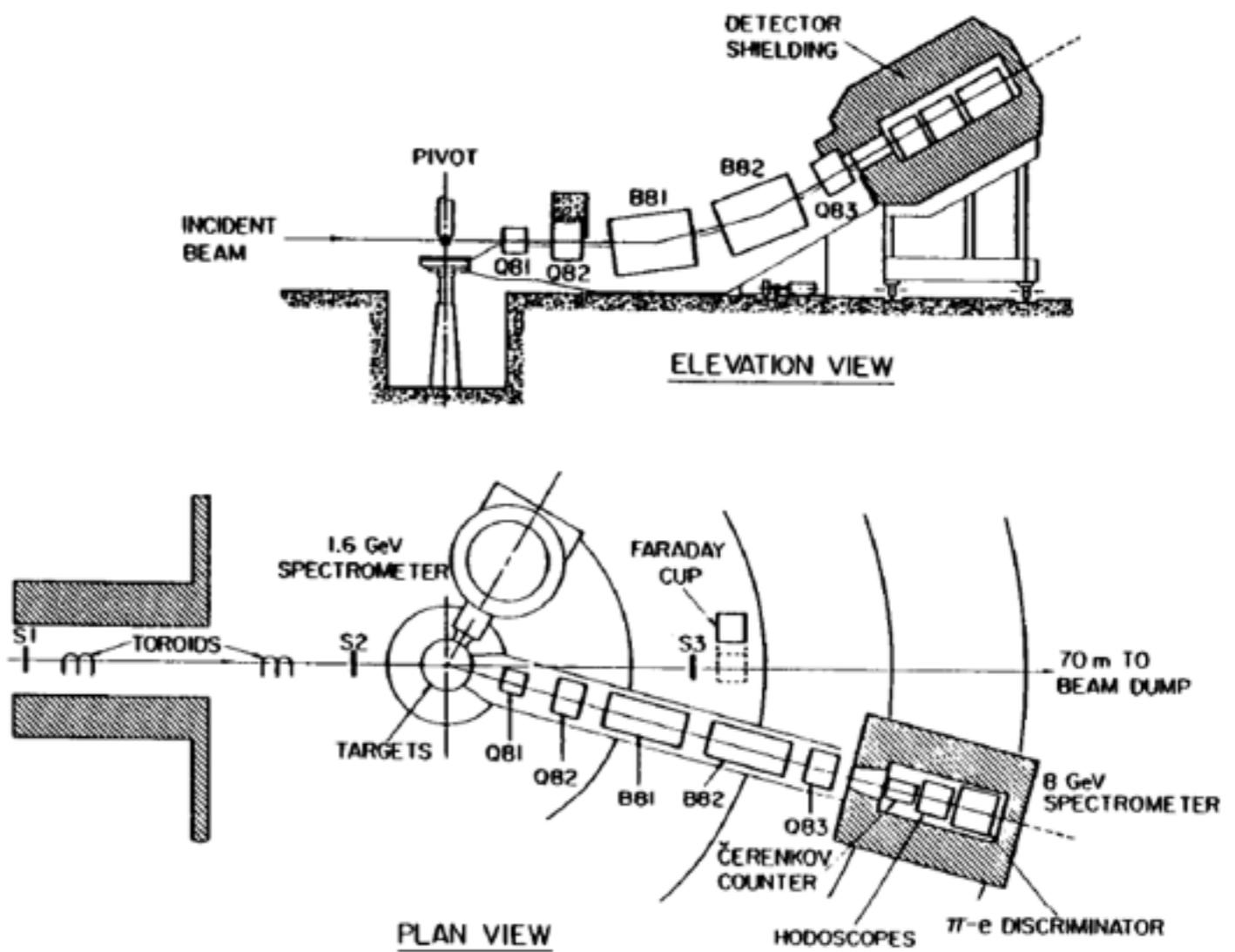


Fig. 1:  $(d^2\sigma/d\Omega dE)/\sigma_{Mott}$ , in  $\text{GeV}^{-1}$ , vs.  $q^2$  for  $W = 2, 3$  and  $3.5$  GeV. The lines drawn through the data are meant to guide the eye. Also shown is the cross section for elastic e-p scattering divided by  $\sigma_{Mott}$ ,  $(d\sigma/d\Omega)/\sigma_{Mott}$ , calculated for  $\theta = 10^\circ$ , using the dipole form factor. The relatively slow variation with  $q^2$  of the inelastic cross section compared with the elastic cross section is clearly shown.



- strong high  $Q^2$  scattering in high energy e-p scattering points to substructure in proton
  - "partons"
  - would later connect to "quarks"

# IGNORANCE

- Form factors are a way to parametrize things we don't know or cannot predict fundamentally.
  - connect it to some physical picture (e.g charge distribution)
  - or appeal to things we do know (Lorentz symmetry) to provide a generalized framework for what it should be
- Typically arise when we have strongly interacting systems
  - e.g. bound quarks in baryons and mesons
    - can't predict a priori all the dynamics, etc.
  - decays of hadrons also have form factors
    - depending on the degrees of freedom, the form factor will have different kinematic dependences.
    - "blob"  $\leftrightarrow$  "form factor"

# NEXT TIME

- Reading
  - 10.1-10.3
  - 10.4
  - 10.5, 10.5.2, 10.6
  - slight departure to original plan, will update.
- Problem Set 2 due next Tuesday