

PHY489/1489

LECTURE 10:
SPIN SUMMATION
ELECTRON-QUARK SCATTERING

SPIN SUMMATION

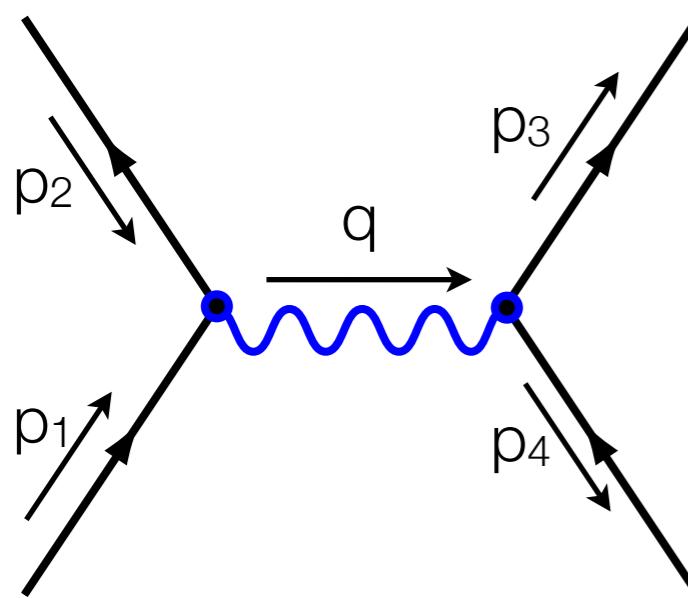
- Last time, we calculated amplitudes using the explicit form of the spinors
 - attractive features: straightforward if tedious math
 - this will always work; it may just be very tedious . . .
- Introduce: spin-summation
 - make use of general spinor properties to derive rather simple expressions
 - preserves Lorentz invariance, easier to deal with masses
 - disadvantage: some new math
- You should have some idea of how both work
 - some things will be (much) easier in one or the other approach.

"COMPLETENESS" RELATIONS

$$\sum_{s=1,2} u^s \bar{u}^s = (\gamma^\mu p_\mu + mc) \quad \sum_{s=1,2} v^s \bar{v}^s = (\gamma^\mu p_\mu - mc)$$

- “outer” product of two Dirac spinors
- sum over spin states s
 - there are two terms in each sum
- Note that these sums are exactly the same as when we:
 - summed over the outgoing spins in a reaction
 - if our detector doesn’t measure the outgoing spin/helicity
 - averaged over incoming spins if the beam is unpolarized
 - this involves summing over all incoming spin configurations and dividing by the total number of configurations
 - e.g. divide by 4 for a pair of incoming unpolarized fermions

AMPLITUDE STRUCTURE



$$\mathcal{M} = -\frac{g_e^2}{(p_1 + p_2)^2} [\bar{u}(3) \gamma^\mu v(4)] [\bar{v}(2) \gamma_\mu u(1)]$$

- need to take the absolute value squared of \mathcal{M}

$$|\mathcal{M}|^2 = \mathcal{M}\mathcal{M}^* = \frac{g_e^4}{(p_1 + p_2)^4} [\bar{u}_3 \gamma^\mu v_4] [\bar{v}_2 \gamma_\mu u_1] \times [\bar{u}_3 \gamma^\nu v_4]^* [\bar{v}_2 \gamma_\nu u_1]^*$$

- Generally, we'll encounter expressions like this:

$$\begin{aligned}
 [\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \Gamma_2 u_b]^* &\quad [\bar{u}_a \Gamma_2 u_b]^* = [\bar{u}_a \Gamma_2 u_b]^\dagger \\
 [\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \Gamma_2 u_b]^* &\quad = [u_a^\dagger \gamma^0 \Gamma_2 u_b]^\dagger = [u_b^\dagger \Gamma_2^\dagger \gamma^{0\dagger} u_a] \\
 [\bar{u}_a \Gamma_1 v_b] [\bar{u}_a \Gamma_2 v_b]^* &\quad = [u_b^\dagger \gamma^0 \gamma^0 \Gamma_2^\dagger \gamma^0 u_a] \\
 [\bar{v}_a \Gamma_1 u_b] [\bar{v}_b \Gamma_2 u_b]^* &\quad = [\bar{u}_b \bar{\Gamma}_2 u_a] = [\bar{u}_b \Gamma_2 u_a]
 \end{aligned}$$

COMPLETENESS RELATION

- We need to calculate terms like this

$$[\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \Gamma_2 u_b]^* = [\bar{u}_a \Gamma_1 u_b] [\bar{u}_b \Gamma_2 u_a]$$

- sum over the spin of b, use the completeness relation

$$\sum_{s_b} [\bar{u}_a \Gamma_2 u_b \bar{u}_b \bar{\Gamma}_2 u_a] \Rightarrow \bar{u}_a \Gamma_2 (\not{p}_b + m_b c) \bar{\Gamma}_2 u_a]$$

$$\sum_{s=1,2} u^s \bar{u}^s = (\gamma^\mu p_\mu + mc)$$

- to sum over the spin of a, write out with indices

$$\sum_i \sum_j [\bar{u}_a]_i [\Gamma_2 (\not{p}_b + m_b c) \bar{\Gamma}_2]_{ij} [u_a]_j$$

- rearrange and introduce spin summation on a

$$\sum_i \sum_j [\Gamma_2 (\not{p}_b + m_b c) \bar{\Gamma}_2]_{ij} \sum_{s_a} [u_a]_j [\bar{u}_a]_i [\not{p}_a + m_a c]_{ji}$$

SUM OVER SPINS OF FINAL STATES

$$\sum_i \sum_j = \left[\Gamma_1(\not{p}_b + m_b c) \bar{\Gamma}_2 \right]_{ij} \left[\not{p}_a + m_a c \right]_{ji}$$

A_{ij}

B_{ji}

sum over j multiplies
the two matrices



$[AB]_{ii}$

$$\sum_i [\Gamma_1(\not{p}_b + m_b c) \bar{\Gamma}_2(\not{p}_a + m_a c)]_{ii}$$

$$\Rightarrow \text{Tr}[\Gamma_1(\not{p}_b + m_b c) \bar{\Gamma}_2(\not{p}_a + m_a c)]$$

$$\sum_{s_a, s_b} [\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \bar{\Gamma}_2 u_b]^* \Rightarrow \text{Tr} [\Gamma_1(\not{p}_b + m_b c) \bar{\Gamma}_2(\not{p}_a + m_a c)]$$

- If particle a or b are antiparticles ("v" spinors), reverse the sign of the mass term

TRACE RELATIONS:

- $\text{Tr}[1] = 4$

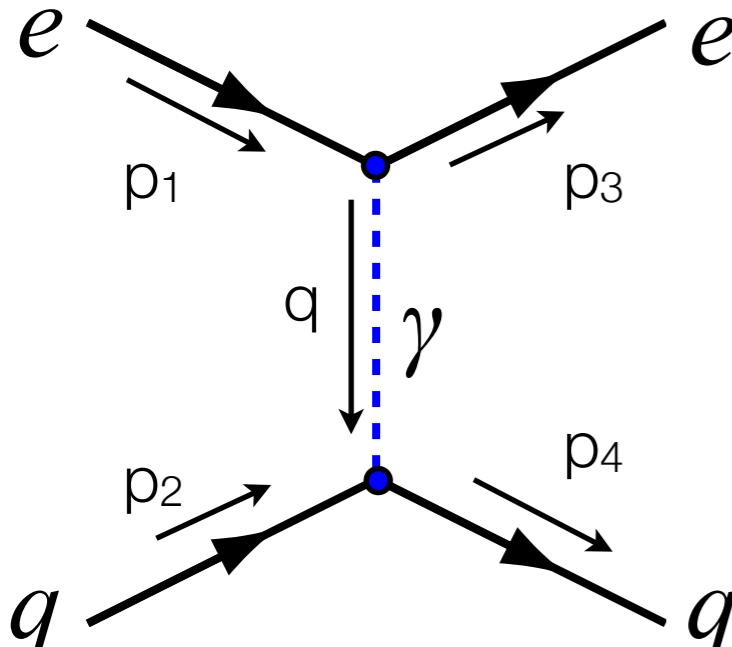
$$\begin{aligned}\text{Tr}(A+B) &= \text{Tr } A + \text{Tr } B \\ \text{Tr}(aA) &= a\text{Tr } A \\ \text{Tr}(AB) &= \text{Tr } (BA)\end{aligned}$$

$$\begin{aligned}\text{Tr}(\gamma^\mu \gamma^\nu) &\quad \text{Tr}(\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu) = \text{Tr}(2g^{\mu\nu}) \\ 2 \times \text{Tr}(\gamma^\mu \gamma^\nu) &= 2 \times \text{Tr}(g^{\mu\nu}) \\ \text{Tr}(\gamma^\mu \gamma^\nu) &= 4 \times g^{\mu\nu}\end{aligned}$$

$$g^{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

- $\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\sigma) = 4 \times (g^{\mu\nu} g^{\lambda\sigma} - g^{\mu\lambda} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\lambda})$
- From here, we have:
 - $\text{Tr}(\not{a} \not{b}) = 4a \cdot b$
 - $\text{Tr}(\not{a} \not{b} \not{c} \not{d}) = 4 [(a \cdot b)(c \cdot d) - (a \cdot c)(b \cdot d) + (a \cdot d)(b \cdot c)]$
- Trace of any single γ matrix is 0
 - $\text{Tr}(\gamma^\mu) = \text{Tr}(\gamma^5 \gamma^5 \gamma^\mu) = -\text{Tr}(\gamma^5 \gamma^\mu \gamma^5) = -\text{Tr}(\gamma^5 \gamma^5 \gamma^\mu)$
 - Likewise, trace of the product of an odd number of γ matrices is 0

ELECTRON-QUARK SCATTERING



$$|\mathcal{M}|^2 = \frac{Q^2 e^4}{(p_1 - p_3)^4} [\bar{u}_3 \gamma^\mu u_1] [\bar{u}_3 \gamma^\nu u_1]^* [\bar{u}_4 \gamma_\mu u_2] [\bar{u}_4 \gamma_\nu u_2]^*$$

$$\begin{aligned} & \bar{u}_3 (-ie\gamma^\mu) u_1 (2\pi)^4 \delta^4(p_1 - p_3 - q) \\ & \bar{u}_4 (-iQe\gamma^\nu) u_2 (2\pi)^4 \delta^4(p_2 - p_4 + q) \\ & \int \frac{dq^4}{(2\pi)^4} \frac{ig_{\mu\nu}}{q^2} \\ & \mathcal{M} = \frac{Qe^2}{(p_1 - p_3)^2} [\bar{u}_3 \gamma^\mu u_1] [\bar{u}_4 \gamma_\mu u_2] \end{aligned}$$

- summing over the sums of all the particles

$$\sum_{s_1, s_3} [\bar{u}_3 \gamma^\mu u_1] [\bar{u}_3 \gamma^\nu u_1]^* \Rightarrow \text{Tr} \left[\gamma^\mu (\not{p}_1 + mc) \gamma^\nu (\not{p}_3 + mc) \right]$$

$$\begin{aligned} & \sum_{s_a, s_b} [\bar{u}_a \Gamma_1 u_b] [\bar{u}_a \bar{\Gamma}_2 u_b]^* \\ & \Rightarrow \text{Tr} \left[\Gamma_1 (\not{p}_b + m_b c) \bar{\Gamma}_2 (\not{p}_a + m_a c) \right] \end{aligned}$$

$$\sum_{s_2, s_4} [\bar{u}_4 \gamma^\mu u_2] [\bar{u}_4 \gamma^\nu u_2]^* \Rightarrow \text{Tr} \left[\gamma_\mu (\not{p}_2 + Mc) \gamma_\nu (\not{p}_4 + Mtc) \right]$$

$$\sum_{\text{spins}} |\mathcal{M}|^2 = \frac{e^4}{(p_1 - p_3)^4} \text{Tr} \left[\gamma^\mu (\not{p}_1 + mc) \gamma^\nu (\not{p}_3 + mc) \right] \text{Tr} \left[\gamma_\mu (\not{p}_2 + Mc) \gamma_\nu (\not{p}_4 + Mc) \right]$$

EVALUATING TRACES

- Here are the trace expressions:

$$\text{Tr} \left[\gamma^\mu (\not{p}_1 + mc) \gamma^\nu (\not{p}_3 + mc) \right] = \text{Tr} \left[\gamma^\mu \not{p}_1 \gamma^\nu \not{p}_3 \right] + m^2 c^2 \text{Tr} [\gamma^\mu \gamma^\nu]$$

$$\text{Tr} \left[\gamma_\mu (\not{p}_2 + Mc) \gamma_\nu (\not{p}_4 + Mc) \right] = \text{Tr} \left[\gamma_\mu \not{p}_2 \gamma_\nu \not{p}_4 \right] + M^2 c^2 \text{Tr} [\gamma^\mu \gamma^\nu]$$

- the relevant trace relations

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\sigma) = 4 \times (g^{\mu\nu} g^{\lambda\sigma} - g^{\mu\lambda} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\lambda})$$

$$\text{Tr}(\gamma^\mu \gamma^\nu) = 4 \times g^{\mu\nu}$$

- we obtain the following:

$$\text{Tr} \left[\gamma^\mu (\not{p}_1 + mc) \gamma^\nu (\not{p}_3 + mc) \right] = 4 \left[p_1^\mu p_3^\nu + p_1^\nu p_3^\mu - (p_1 \cdot p_3) g^{\mu\nu} + m^2 c^2 g^{\mu\nu} \right]$$

$$\text{Tr} \left[\gamma_\mu (\not{p}_2 + Mc) \gamma_\nu (\not{p}_4 + Mc) \right] = 4 \left[p_{2\mu} p_{4\nu} + p_{2\nu} p_{4\mu} - (p_2 \cdot p_4) g_{\mu\nu} + M^2 c^2 g^{\mu\nu} \right]$$

END GAME

- Contract the two trace expressions

$$4 [p_1^\mu p_3^\nu + p_1^\nu p_3^\mu - (p_1 \cdot p_3)g^{\mu\nu} + m^2 c^2 g^{\mu\nu}]$$

$$4 [p_{2\mu} p_{4\nu} + p_{2\nu} p_{4\mu} - (p_2 \cdot p_4)g_{\mu\nu} + M^2 c^2 g_{\mu\nu}]$$

$$\Rightarrow 32 [(p_1 \cdot p_2)(p_3 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) - (M^2 c^2)(p_1 \cdot p_3) - m^2 c^2 (p_2 \cdot p_4) + 2m^2 c^2 M^2 c^2]$$

- Putting it back into the amplitude expression

$$\sum_{\text{spins}} |\mathcal{M}|^2 = \frac{32Q^2 e^4}{(p_1 - p_3)^4} [(p_1 \cdot p_2)(p_3 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) - (M^2 c^2)(p_1 \cdot p_3) - m^2 c^2 (p_2 \cdot p_4) + 2m^2 c^2 M^2 c^2]$$

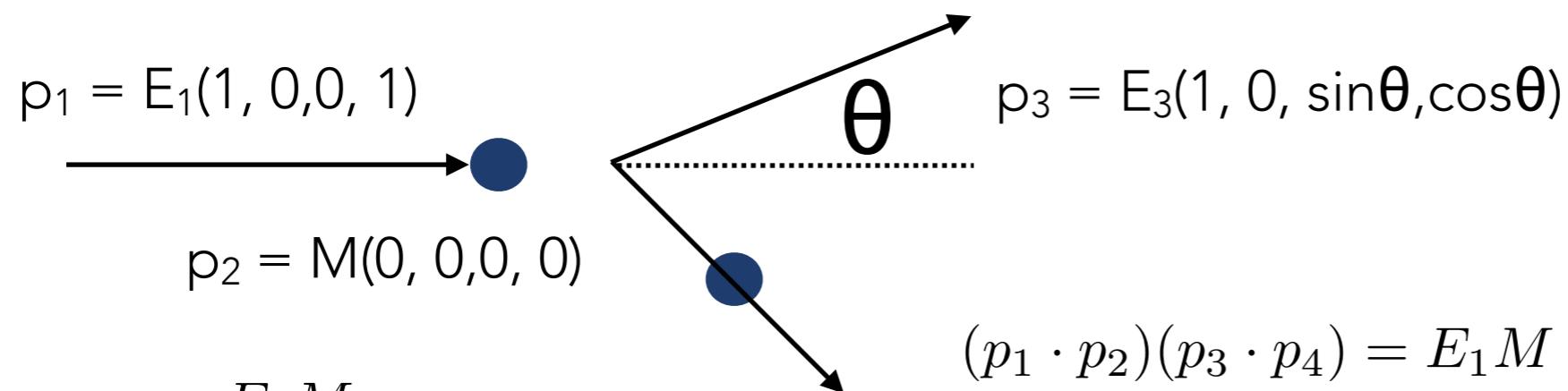
- don't forget we need 1/4 averaging factor
- Put it into the phase space expression for lab-frame scattering (note this expression ignores the electron mass)

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left(\frac{E_3}{ME_1} \right)^2 |\mathcal{M}|^2$$

KINEMATICS:

- First let's drop the electron mass

$$\langle |\mathcal{M}|^2 \rangle = \frac{8Q^2 e^4}{(p_1 - p_3)^4} [(p_1 \cdot p_2)(p_3 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) - (M^2 c^2)(p_1 \cdot p_3)]$$



$$p_1 \cdot p_2 = E_1 M$$

$$p_2 \cdot p_3 = E_3 M$$

$$p_1 \cdot p_3 = E_1 E_3 (1 - \cos \theta)$$

$$p_1 \cdot p_4 = p_1 \cdot (p_1 + p_2 - p_3) \sim p_1 \cdot p_2 - p_1 \cdot p_3 = E_1 M - E_1 E_3 (1 - \cos \theta)$$

$$p_3 \cdot p_4 = p_3 \cdot (p_1 + p_2 - p_3) \sim p_1 \cdot p_3 - p_2 \cdot p_3 = E_1 E_3 (1 - \cos \theta) + E_3 M$$

$$(p_1 \cdot p_2)(p_3 \cdot p_4) = E_1 M \times (E_1 E_3 (1 - \cos \theta) + E_3 M)$$

$$(p_1 \cdot p_4)(p_2 \cdot p_3) = E_3 M \times (E_1 M - E_1 E_3 (1 - \cos \theta))$$

$$M^2 c^2 (p_1 \cdot p_3) = M^2 E_1 E_3 (1 - \cos \theta)$$

$$\langle |\mathcal{M}|^2 \rangle = \frac{8Q^2 e^4}{(p_1 - p_3)^4} E_1 E_3 M [(E_1 - E_3)(1 - \cos \theta) + M(1 + \cos \theta)]$$

$$= \frac{16Q^2 e^4}{(p_1 - p_3)^4} E_1 E_3 M \left[(E_1 - E_3) \sin^2 \frac{\theta}{2} + M \cos^2 \frac{\theta}{2} \right]$$

MOMENTUM TRANSFER:

$$\langle |\mathcal{M}|^2 \rangle = \frac{16Q^2e^4}{(p_1 - p_3)^4} E_1 E_3 M \left[(E_1 - E_3) \sin^2 \frac{\theta}{2} + M \cos^2 \frac{\theta}{2} \right]$$

$$(p_1 - p_3)^2 \sim -2p_1 \cdot p_3 = -4E_1 E_3 \sin^2 \frac{\theta}{2} = q^2 \equiv -Q^2$$

$$Q^2 = 2M(E_1 - E_3)$$

$$\langle |\mathcal{M}|^2 \rangle = \frac{M^2 e^4}{E_1 E_3 \sin^4 \frac{\theta}{2}} \left[\cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right]$$

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left(\frac{E_3}{ME_1} \right)^2 |\mathcal{M}|^2 \quad e^2 = 4\pi\alpha$$

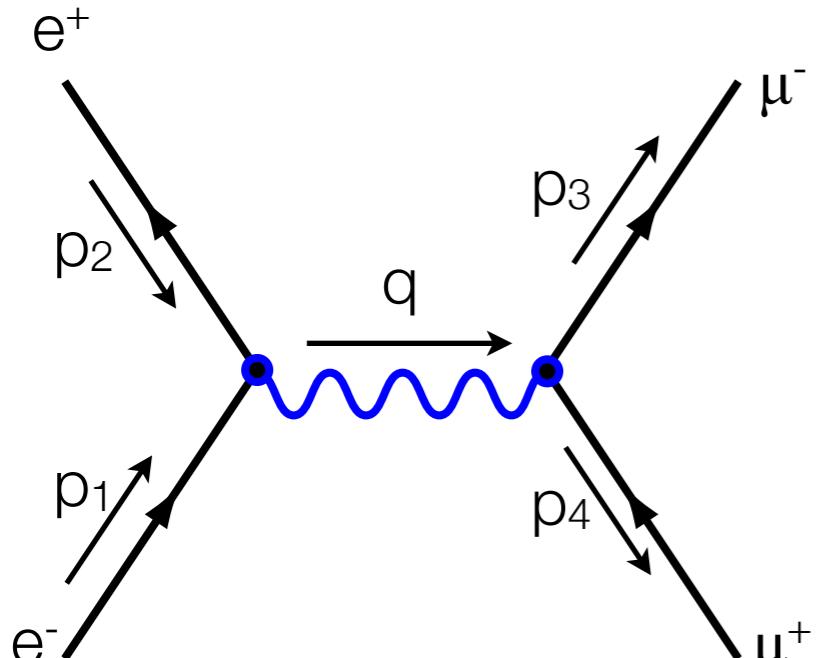
$$= \frac{e^4}{64\pi^2} \frac{1}{E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[\cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right] = \frac{\alpha^2}{4} \frac{1}{E_1^2 \sin^4 \frac{\theta}{2}} \frac{E_3}{E_1} \left[\cos^2 \frac{\theta}{2} + \frac{Q^2}{2M} \sin^2 \frac{\theta}{2} \right]$$

A FEW NOTES

- Note that there is only one degree of freedom in the scattering:
 - $E_3 = \frac{E_1 M}{M + E_1(1 - \cos \theta)}$
 - $Q^2 = \frac{2ME_1^2(1 - \cos \theta)}{M + E_1(1 - \cos \theta)}$
 - i.e. if we know one of E_3 , Q , or Q^2 , the others are determined.

NEXT TIME:

- Please read
 - 7.5, 9.1
 - if you have a chance, 10.1



$$\mathcal{M} = -\frac{g_e^2}{(p_1 + p_2)^2} [\bar{u}(3) \gamma^\mu v(4)] [\bar{v}(2) \gamma_\mu u(1)]$$

$$\mathcal{M}^* = -\frac{g_e^2}{(p_1 + p_2)^2} [\bar{u}(3) \gamma^\nu v(4)]^* [\bar{v}(2) \gamma_\nu u(1)]^*$$

$$|\mathcal{M}|^2 = \frac{g_e^4}{(p_1 + p_2)^4} [\bar{u}(3) \gamma^\mu v(4)] [\bar{u}(3) \gamma^\nu v(4)]^* [\bar{v}(2) \gamma_\mu u(1)] [\bar{v}(2) \gamma_\nu u(1)]^*$$

$$\sum_{\text{spins}} [\bar{u}(3) \gamma^\mu v(4)] [\bar{u}(3) \gamma^\nu v(4)]^* = \text{Tr} [(\gamma^\mu (\not{p}_4 - Mc) \gamma^\nu (\not{p}_3 + Mc))]$$

$$\sum_{\text{spins}} [\bar{v}(2) \gamma^\mu u(1)] [\bar{v}(2) \gamma^\nu u(1)]^* = \text{Tr} [\gamma_\mu (\not{p}_1 + mc) \gamma_\nu (\not{p}_2 - mc)]$$

$$\sum_{\text{spins}} [\bar{u}(3)\gamma^\mu v(4)] [\bar{u}(3)\gamma^\nu v(4)]^* = \text{Tr} [(\gamma^\mu(\not{p}_4 - Mc)\gamma^\nu(\not{p}_3 + Mc)]$$

$$\sum_{\text{spins}} [\bar{v}(2)\gamma^\mu u(1)] [\bar{v}(2)\gamma^\nu u(1)]^* = \text{Tr} [\gamma_\mu(\not{p}_1 + mc)\gamma_\nu(\not{p}_2 - mc)]$$

$$\text{Tr} \left[\gamma^\mu(\not{p}_4 - Mc)\gamma^\nu(\not{p}_3 + Mc) \right] = 4 \left(p_4^\mu p_3^\nu + p_4^\nu p_3^\mu - g^{\mu\nu}(p_3 \cdot p_4) - 4Mc^2 g^{\mu\nu} \right)$$

$$\text{Tr} \left[\gamma_\mu(\not{p}_1 + mc)\gamma_\nu(\not{p}_2 - mc) \right] = 4 \left(p_{1\mu} p_{2\nu} + p_{1\nu} p_{2\mu} - g_{\mu\nu}(p_1 \cdot p_2) - 4mc^2 g_{\mu\nu} \right)$$