

The Z boson

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Final examination

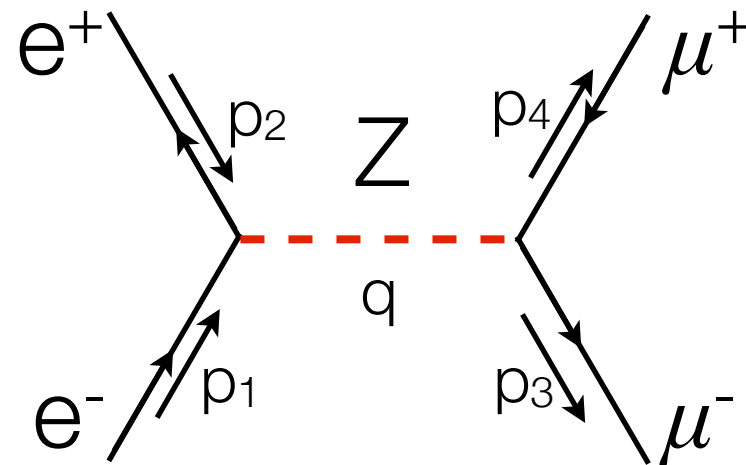
- Feynman diagram calculation:
 - electromagnetic or weak interaction process (scattering/decay)
 - use Feynman rules and get to a cross section or decay rate
 - there are many steps, so I will try to provide the result after each step
 - if you get stuck on any step, take the result and go to the next step
- Structure of the weak interaction
 - parity violation and its consequences (e.g. what decays are suppressed)
 - CKM matrix elements (how to apply them)
 - electroweak unification:
 - what is it?
 - what is θ_w ? where does it appear?
- You can use your textbook and a non-programmable calculator.

Averaging with projection operators

- suppose we use P_L to project out the left chiral component of an incoming particle. The other incoming particle is unpolarized.
 - In summing the spins, we get: $|M_{LL}|^2 + |M_{LR}|^2 + |M_{RL}|^2 + |M_{RR}|^2$
 - The projection operator will kill $|M_{RL}|^2$ and $|M_{RR}|^2$, leaving $|M_{LL}|^2 + |M_{LR}|^2$
 - The average amplitude is $1/2 \times (|M_{LL}|^2 + |M_{LR}|^2)$
 - so we divide by 2
- If we use P_L to project out the left chiral components of both particles:
 - In summing the spins, we get: $|M_{LL}|^2 + |M_{LR}|^2 + |M_{RL}|^2 + |M_{RR}|^2$
 - The projection operator will kill $|M_{RL}|^2$ and $|M_{RR}|^2$, leaving $|M_{LL}|^2 + |M_{LR}|^2$
 - only leaving $|M_{LL}|^2$
 - The “average” amplitude is $|M_{LL}|^2$
 - so we leave it alone.

$$e^+ e^- \rightarrow \mu^+ \mu^-$$

- Let's consider the pair production of muons via the Z:



Note: different convention from text book (p_3 assigned here to μ^- , not μ^+)

Assume e, μ masses can be ignored

- Feynman rules:

- incoming leg: $\bar{v}(2) \frac{-ig_Z}{2} \gamma^\mu (c_V^e - c_A^e \gamma^5) u(1) \quad (2\pi)^4 \delta^4(p_1 + p_2 - q)$

- outgoing leg: $\bar{u}(3) \frac{-ig_Z}{2} \gamma^\nu (c_V^\mu - c_A^\mu \gamma^5) v(4) \quad (2\pi)^4 \delta^4(q - p_3 + p_4)$

- Propagator: $\int \frac{d^4 q}{(2\pi)^4} \frac{-i(g_{\mu\nu} - q_\mu q_\nu / m_Z^2 c^2)}{q^2 - m_Z^2 c^2}$

$$\mathcal{M} = -\frac{g_Z^2}{4(q^2 - m_Z^2 c^2)} [\bar{v}(2) \gamma^\mu (c_V^e - c_A^e \gamma^5) u(1)] [\bar{u}(3) \gamma^\nu (c_V^\mu - c_A^\mu \gamma^5) v(4)]$$

$$[g_{\mu\nu} - q_\mu q_\nu / (M_z^2 c^2)]$$

Spin Summation

$$\mathcal{M} = -\frac{g_Z^2}{4(q^2 - m_Z^2 c^2)} [\bar{v}(2)\gamma^\mu(c_V^e - c_A^e\gamma^5)u(1)] [\bar{u}(3)\gamma_\mu(c_V^\mu - c_A^\mu\gamma^5)v(4)]$$

$$\mathcal{M}\mathcal{M}^* = \left[\frac{g_Z^2}{4(q^2 - m_Z^2 c^2)} \right]^2 \frac{[\bar{v}(2)\gamma^\mu(c_V^e - c_A^e\gamma^5)u(1)] [\bar{v}(2)\gamma^\nu(c_V^e - c_A^e\gamma^5)u(1)]^*}{[\bar{u}(3)\gamma_\mu(c_V^\mu - c_A^\mu\gamma^5)v(4)] [\bar{u}(3)\gamma_\nu(c_V^\mu - c_A^\mu\gamma^5)v(4)]^*}$$

$$\sum_{a, b \text{ spins}} [\bar{u}(a)\Gamma_1 u(b)] [\bar{u}(a)\bar{\Gamma}_2 u(b)]^* = \text{Tr} [\Gamma_1(\not{p}_b + m_b c)\bar{\Gamma}_2(\not{p}_a + m_a c)]$$

$$\langle \mathcal{M}\mathcal{M}^* \rangle = \frac{1}{4} \left[\frac{g_Z^2}{4(q^2 - m_Z^2 c^2)} \right]^2 \frac{\text{Tr} [\gamma^\mu(c_V^e - c_A^e\gamma^5) \not{p}_1 \gamma^\nu(c_V^e - c_A^e\gamma^5) \not{p}_2]}{\text{Tr} [\gamma_\mu(c_V^\mu - c_A^\mu\gamma^5) \not{p}_4 \gamma_\nu(c_V^\mu - c_A^\mu\gamma^5) \not{p}_3]}$$

Breaking up the Terms:

- The first trace:

$$\text{Tr} [\gamma^\mu (c_V^e - c_A^e \gamma^5) \not{p}_1 \gamma^\nu (c_V^e - c_A^e \gamma^5) \not{p}_2]$$

$$\text{Tr} [\gamma^\mu \not{p}_1 \gamma^\nu \not{p}_2] \Rightarrow 4 [p_1^\mu p_2^\nu + p_1^\nu p_2^\mu - g^{\mu\nu} (p_1 \cdot p_2)]$$

$$\text{Tr} [\gamma^5 \not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu] \Rightarrow 4i\epsilon^{\alpha\nu\beta\mu} p_{1\mu} p_{2\nu}$$

$$4(c_V^{e2} + c_A^{e2}) [p_1^\mu p_2^\nu + p_1^\nu p_2^\mu - g^{\mu\nu} (p_1 \cdot p_2)] - 8ic_V^e c_A^e \times \epsilon^{\alpha\nu\beta\mu} p_{1\mu} p_{2\nu}$$

- Second trace:

$$\text{Tr} [\gamma_\mu (c_V^\mu - c_A^\mu \gamma^5) \not{p}_4 \gamma_\nu (c_V^e - c_A^e \gamma^5) \not{p}_3]$$

$$4(c_V^{\mu2} + c_A^{\mu2}) [p_{3\mu} p_{4\nu} + p_{3\nu} p_{4\mu} - g_{\mu\nu} (p_3 \cdot p_4)] - 8ic_V^\mu c_A^\mu \times \epsilon_{\rho\nu\sigma\mu} p_3^\rho p_4^\sigma$$

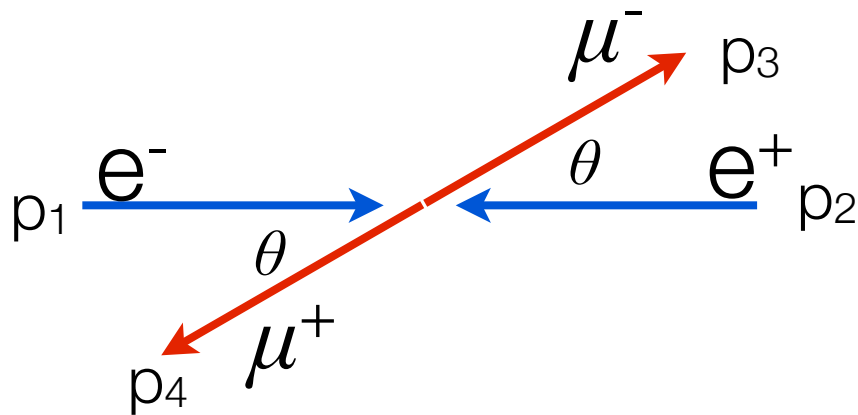
Our final expression for the amplitude:

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{2} \left[\frac{g_Z^2}{(q^2 - m_Z^2 c^2)} \right]^2$$

$$\{ (c_V^e)^2 + (c_A^e)^2 \} (c_V^\mu)^2 + (c_A^\mu)^2 \times [(p_1 \cdot p_4)(p_2 \cdot p_3) + (p_1 \cdot p_3)(p_2 \cdot p_4)]$$

$$+ 4c_V^e c_A^e c_V^\mu c_A^\mu [(p_1 \cdot p_4)(p_2 \cdot p_3) - (p_1 \cdot p_3)(p_2 \cdot p_4)] \}$$

- Let's put in the kinematics: $q^2 = (p_1 + p_2)^2 = p_1^2 + p_2^2 + 2p_1 \cdot p_2 \Rightarrow 4E^2$



$$p_1 \cdot p_3 = E^2/c^2 - \mathbf{p}^2 \cos \theta \Rightarrow E^2(1 - \cos \theta)$$

$$p_1 \cdot p_4 = E^2/c^2 + \mathbf{p}^2 \cos \theta \Rightarrow E^2(1 + \cos \theta)$$

$$p_2 \cdot p_3 = E^2/c^2 + \mathbf{p}^2 \cos \theta \Rightarrow E^2(1 + \cos \theta)$$

$$p_2 \cdot p_4 = E^2/c^2 - \mathbf{p}^2 \cos \theta \Rightarrow E^2(1 - \cos \theta)$$

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{2} \left[\frac{g_Z^2 E^2}{(4E^2 - m_Z^2 c^2)} \right]^2$$

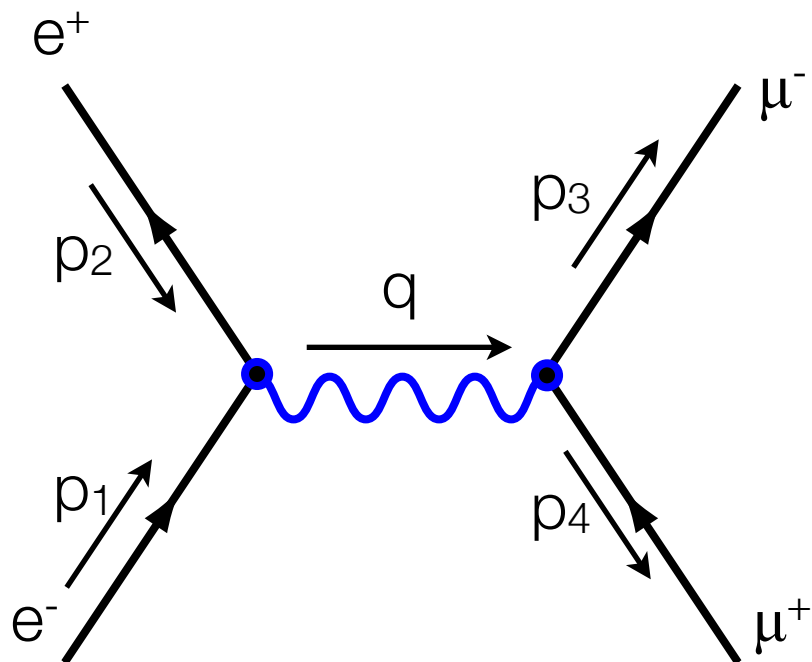
$$\{ [(c_V^e)^2 + (c_A^e)^2] [(c_V^\mu)^2 + (c_A^\mu)^2] \times [2(1 + \cos^2 \theta)] + 16 c_V^e c_A^e c_V^\mu c_A^\mu \cos \theta \}$$

The End:

- The differential cross section in the CM frame

$$\frac{d\sigma}{d\Omega} = \left(\frac{\hbar c}{8\pi} \right)^2 \frac{S |\mathcal{M}|^2}{(E_1 + E_2)^2} \frac{|p_f|}{|p_i|}$$

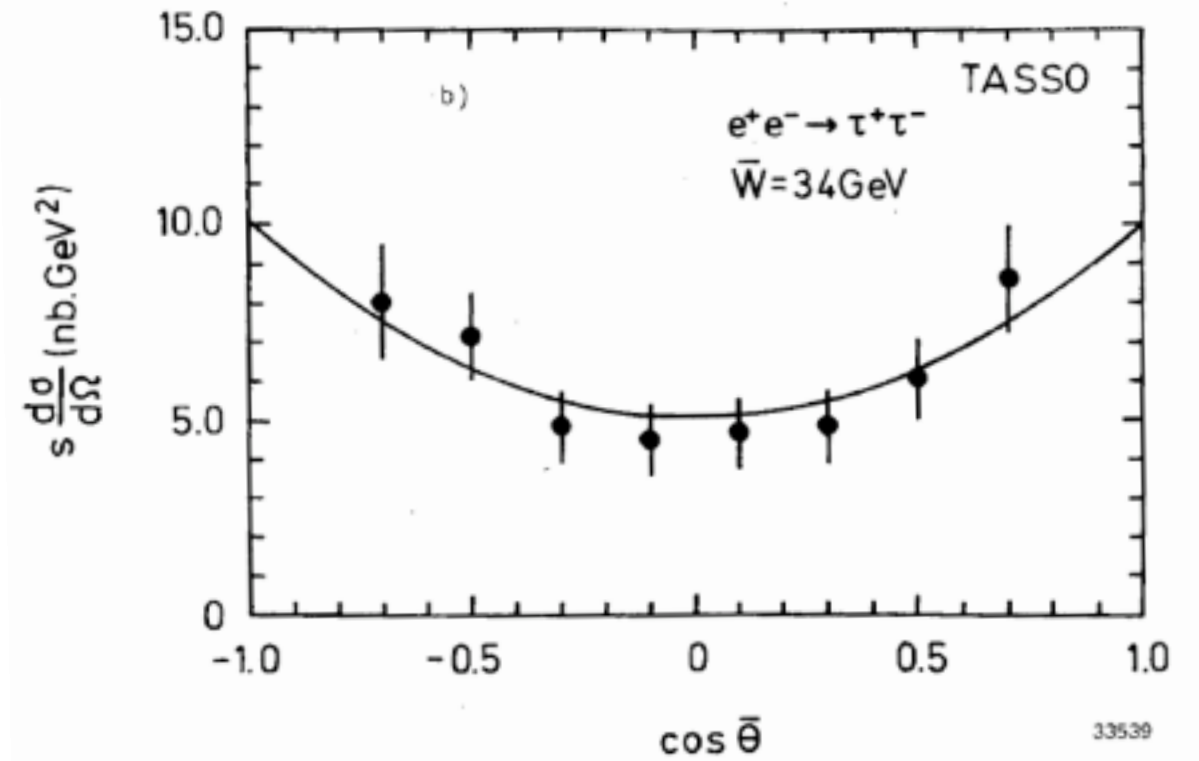
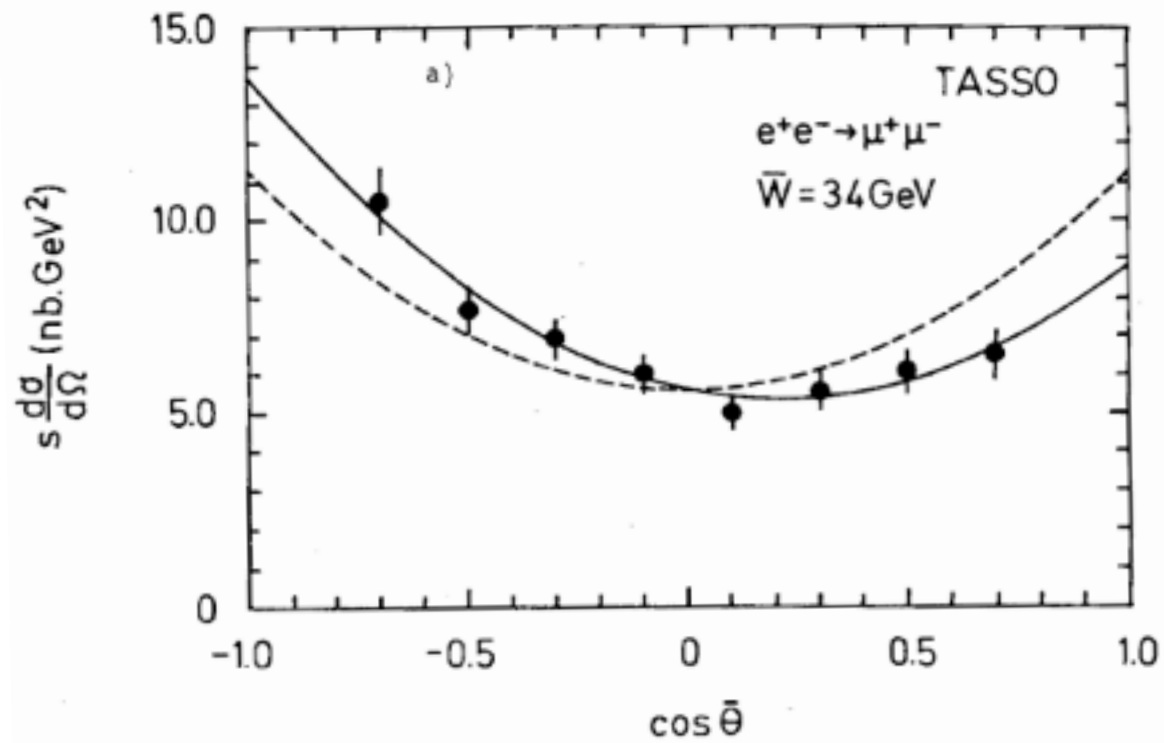
$$\frac{d\sigma}{d\Omega} = \left(\frac{\hbar c}{16\pi} \right)^2 \left[\frac{g_Z^2 E}{(4E^2 - m_Z^2 c^2)} \right]^2 \{ [(c_V^e)^2 + (c_A^e)^2] [(c_V^\mu)^2 + (c_A^\mu)^2] \times (1 + \cos^2 \theta) + 8 c_V^e c_A^e c_V^\mu c_A^\mu \cos \theta \}$$



Compare to electrodynamics

$$\frac{d\sigma}{d\Omega} = \left(\frac{\hbar c}{8\pi} \right)^2 \frac{g_e^4}{4E^2} [1 + \cos^2 \theta]$$

Comparison to measurements



How did the asymmetry come about?

- We can break up the Z interaction into two parts:

$$c_V \gamma^\mu + c_A \gamma^\mu \gamma^5 \Rightarrow \frac{1}{2}(c_V + c_A) \gamma^\mu (1 + \gamma^5) + \frac{1}{2}(c_V - c_A) \gamma^\mu (1 - \gamma^5)$$

- One couples to left chiral and the other to right chiral particles
 - But the coupling is not the same!
- Within the incident e^+e^- beams, we have assumed there is no polarization
 - equal components of each chirality
- Recall from last time that:

$$\gamma^\mu \left(\frac{1 - \gamma^5}{2} \right) = \left(\frac{1 + \gamma^5}{2} \right) \gamma^\mu \left(\frac{1 - \gamma^5}{2} \right)$$

- First term above couples right chiral electrons with left chiral positrons

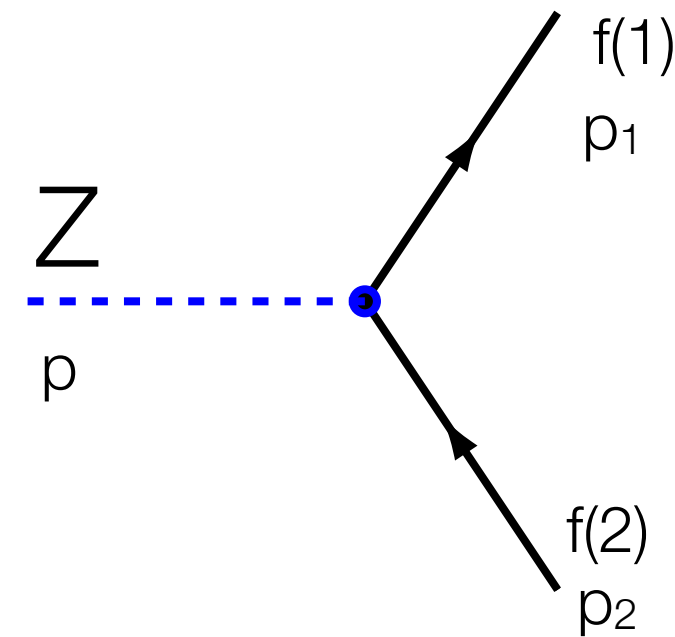
$$u_L \equiv \frac{1}{2}(1 - \gamma^5)u \quad \bar{u}_R \equiv \bar{u} \frac{1}{2}(1 + \gamma^5)$$

- vice versa for the second term

Z has net polarization

Z decays: (9.23)

- Consider the decay of the Z to two fermions (quark/lepton)
 - let's ignore the fermion masses (i.e. set to 0)
- Treat all possible fermion pairs by not specifying c_V , c_A
 - putting in the appropriate values at the end, case-by-case



Vertex $\bar{u}_1 \frac{-ig_Z}{2} \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2 \times (2\pi)^4 \delta^4(p - p_1 - p_2)$

Incoming Z ϵ_μ

$$\mathcal{M} = \frac{g_Z}{2} \epsilon_\mu \times \bar{u}_1 \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2$$

$$|\mathcal{M}|^2 = \frac{g_Z^2}{4} \epsilon_\mu \epsilon_\nu^* \times [\bar{u}_1 \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2] [\bar{u}_1 \gamma^\nu (c_V^f - c_A^f \gamma^5) v_2]^*$$

Spin/polarization summation

$$|\mathcal{M}|^2 = \frac{g_Z^2}{4} \epsilon_\mu \epsilon_\nu^* \times [\bar{u}_1 \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2] [\bar{u}_1 \gamma^\nu (c_V^f - c_A^f \gamma^5) v_2]^*$$

- Hopefully the fermion part is familiar to you now

$$\sum_{s_1, s_2} [\bar{u}_1 \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2] [\bar{u}_1 \gamma^\mu (c_V^f - c_A^f \gamma^5) v_2]^*$$

$$\Rightarrow \text{Tr}[\not{p}_2 \gamma^\mu (c_V^f - c_A^f \gamma^5) \not{p}_1 \gamma^\nu (c_V^f - c_A^f \gamma^5)]$$

- what about the Z?

$$\sum_{s_Z} \epsilon_\mu \epsilon_\nu^* \Rightarrow -g_{\mu\nu} + \frac{p_\mu p_\nu}{M_Z^2 c^2} \quad \text{from Problem 9.1 in PS 4}$$

Trace algebra

$$\text{Tr}[\not{p}_2 \gamma^\mu (c_V^f - c_A^f \gamma^5) \not{p}_1 \gamma^\nu (c_V^f - c_A^f \gamma^5)]$$

- As we discussed just before, break this up into parts we know how to manage

$$\begin{aligned}\text{Tr}[\not{p}_2 \gamma^\mu \not{p}_1 \gamma^\nu] &= 4 \times [p_2^\mu p_1^\nu + p_2^\nu p_1^\mu - g^{\mu\nu} (p_1 \cdot p_2)] \\ &\Rightarrow 4(c_V^f{}^2 + c_A^f{}^2) \times [p_2^\mu p_1^\nu + p_2^\nu p_1^\mu - g^{\mu\nu} (p_1 \cdot p_2)]\end{aligned}$$

$$\begin{aligned}\text{Tr}[\gamma^5 \not{p}_1 \gamma^\nu \not{p}_2 \gamma^\mu] &= 4i \epsilon^{\alpha\nu\beta\mu} p_{1\alpha} p_{2\beta} \\ &\Rightarrow -8i c_V^f c_A^f \times \epsilon^{\alpha\nu\beta\mu} p_{1\alpha} p_{2\beta} \\ &\Rightarrow -8i c_V^f c_A^f \times \epsilon^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta}\end{aligned}$$

$$\Rightarrow 4(c_V^2 + c_A^2) \times [p_2^\mu p_1^\nu + p_2^\nu p_1^\mu - g^{\mu\nu} (p_1 \cdot p_2)] - 8i c_V^f c_A^f \times \epsilon^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta}$$

Put it together:

$$\begin{aligned}
 |\mathcal{M}|^2 &= g_Z^2 \left[-g_{\mu\nu} + \frac{p_\mu p_\nu}{M_Z^2 c^2} \right] \\
 &\quad \times [(c_V^2 + c_A^2) \times [p_2^\mu p_1^\nu + p_2^\nu p_1^\mu - g^{\mu\nu} (p_1 \cdot p_2)] \\
 &\quad - 2i c_V^f c_A^f \times \epsilon^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta}]
 \end{aligned}$$

$$\begin{aligned}
 |\mathcal{M}|^2 &= g_Z^2 (c_V^2 + c_A^2) \times \left[-(p_1 \cdot p_2) - (p_1 \cdot p_2) + 4(p_1 \cdot p_2) \right. \\
 &\quad \left. + \frac{1}{m_Z^2 c^2} ((p \cdot p_2)(p \cdot p_1) + (p \cdot p_2)(p \cdot p_1) - p^2 (p_1 \cdot p_2)) \right]
 \end{aligned}$$

$$|\mathcal{M}|^2 = g_Z^2 (c_V^2 + c_A^2) \times \left[2(p_1 \cdot p_2) + \frac{1}{m_Z^2 c^2} ((2(p \cdot p_2)(p \cdot p_1) - p^2 (p_1 \cdot p_2)) \right]$$

$$p^2 = (p_1 + p_2)^2 = p_1^2 + p_2^2 + 2(p_1 \cdot p_2) = 2(p_1 \cdot p_2) = m_Z^2 c^2$$

$$p \cdot p_{1,2} = (p_1 + p_2) \cdot p_{1,2} = p_{1,2}^2 + (p_1 \cdot p_2)$$

$$|\mathcal{M}|^2 = g_Z^2 (c_V^2 + c_A^2) m_Z^2 c^2$$

Towards a decay rate:

- Averaging over initial spins:

$$|\mathcal{M}|^2 = g_Z^2 (c_V^2 + c_A^2) m_Z^2 c^2$$

$$\langle |\mathcal{M}|^2 \rangle \Rightarrow \frac{1}{3} g_Z^2 (c_V^2 + c_A^2) m_Z^2 c^2$$

$$\Gamma = \frac{S |\mathbf{p}|}{8\pi \hbar m_Z^2 c} |\mathcal{M}|^2$$

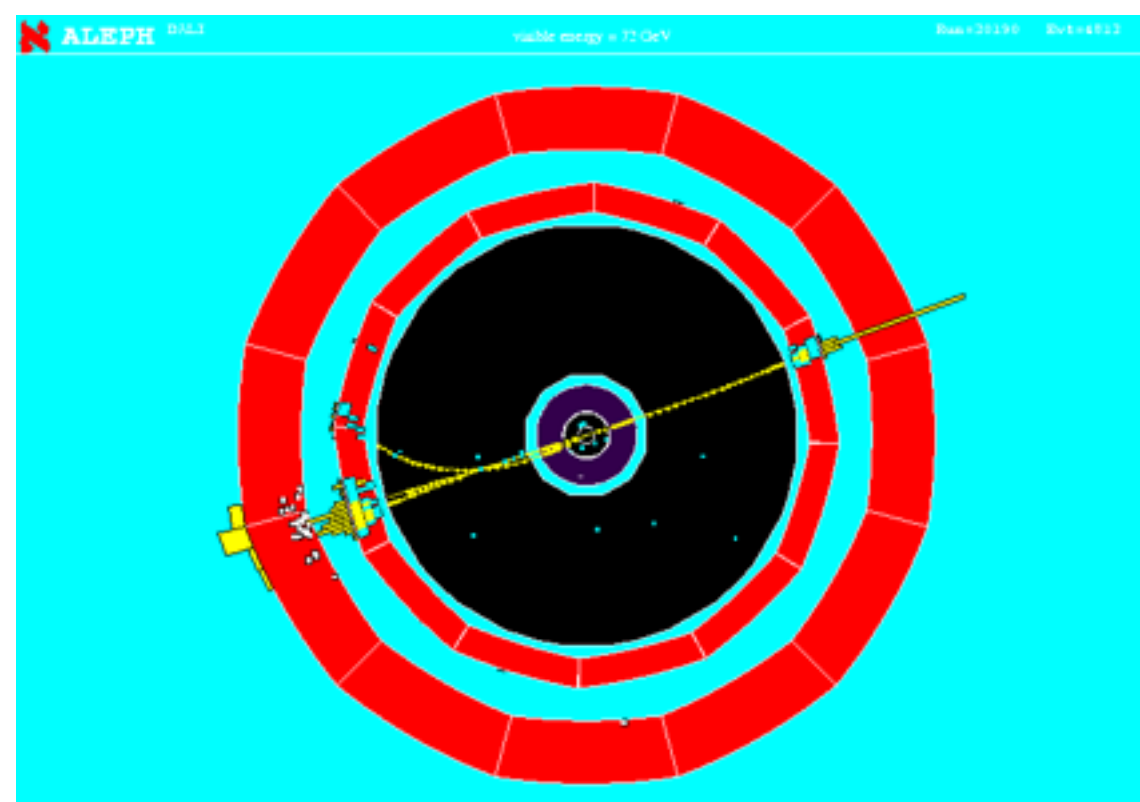
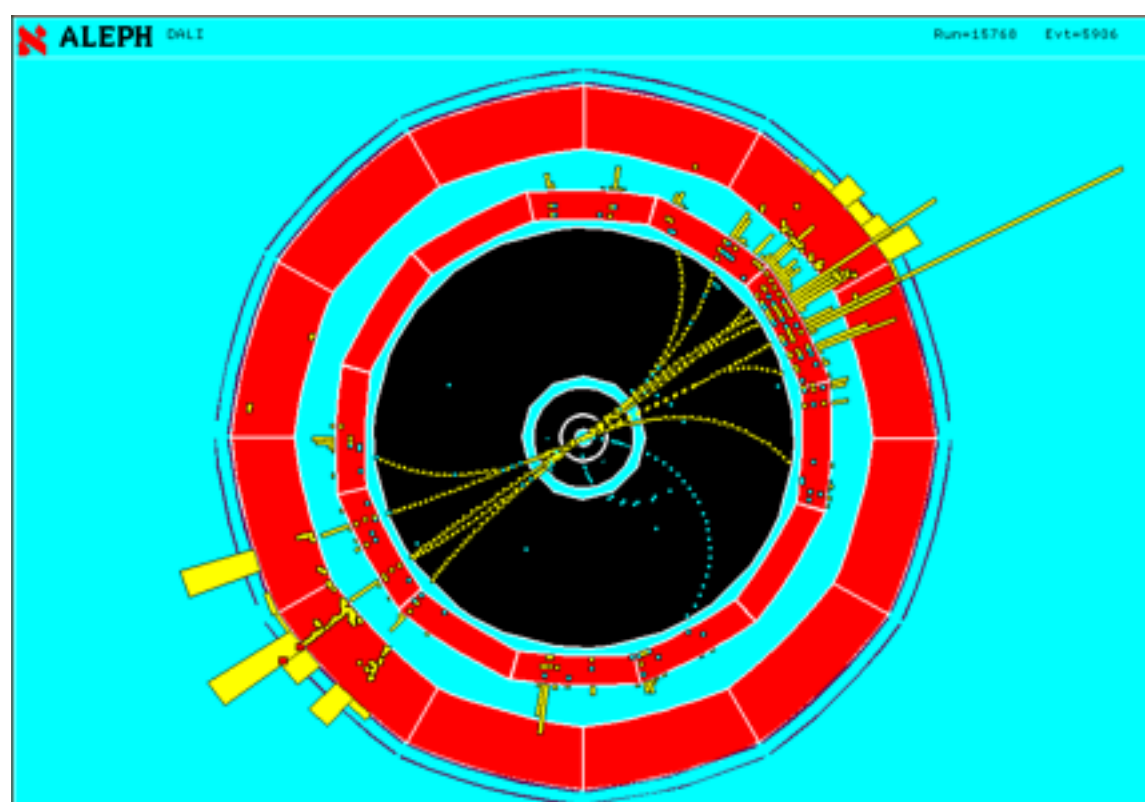
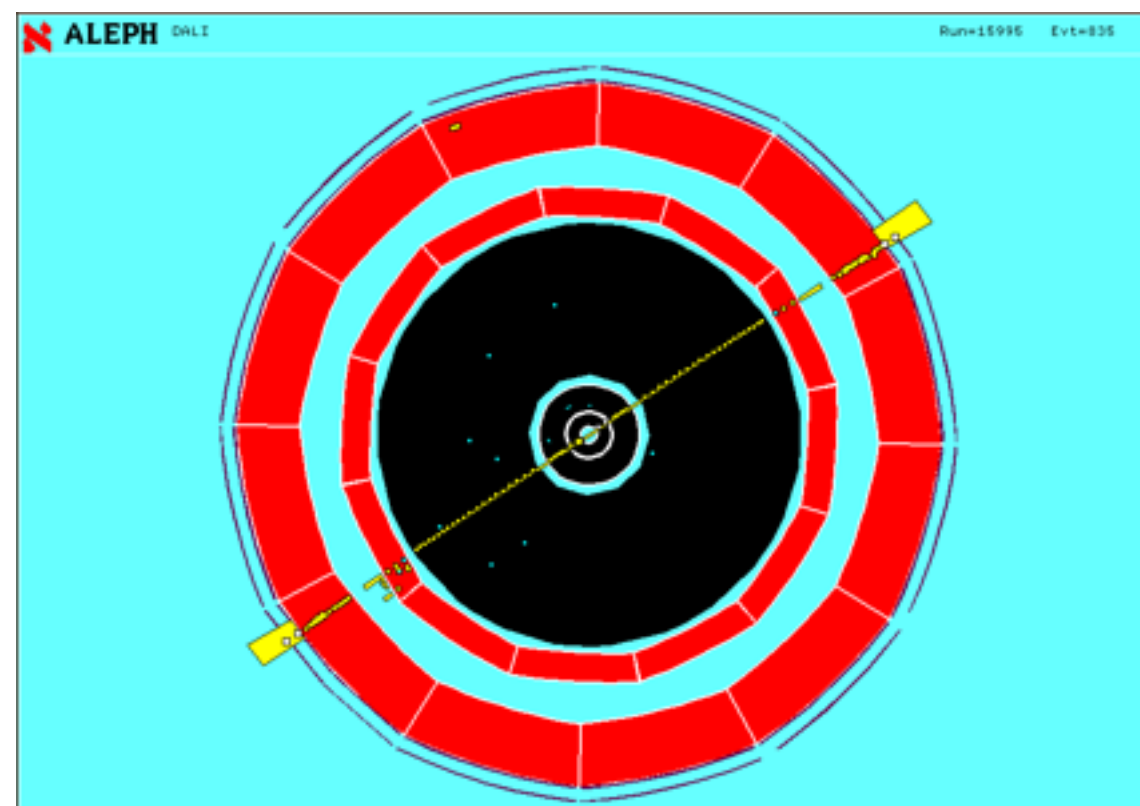
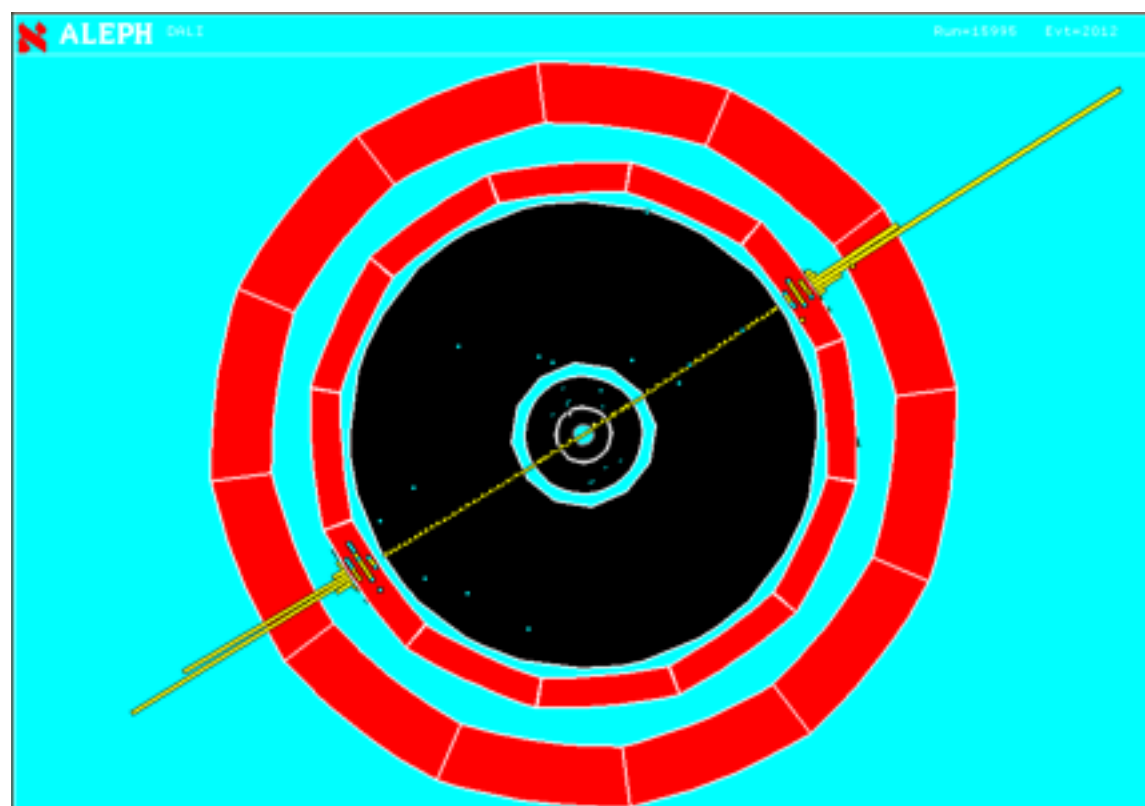
$$|\mathbf{p}| = m_Z c / 2$$

$$S = 1$$

$$\Rightarrow \frac{1}{3} g_Z^2 (c_V^2 + c_A^2) m_Z^2 c^2 \frac{m_Z c}{2 \times 8\pi m_Z^2 c}$$

$$\Rightarrow \frac{1}{48\pi \hbar} g_Z^2 (c_V^2 + c_A^2) m_Z c^2$$

Z decays at LEP



Prediction vs. Measurement

	C_V	C_A	$C_V^2 + C_A^2$	Fraction
$\nu_e \nu_\mu \nu_\tau$	1/2	1/2	0.5 (x3)	6.8%
$e \mu \tau$	$-1/2 + 2 \sin^2 \theta_W$	-1/2	0.251 (x3)	3.4%
$u c t$	$+1/2 - 4/3 \sin^2 \theta_W$	1/2	0.287 (x2x3)	11.6%
$d s b$	$-1/2 + 2/3 \sin^2 \theta_W$	-1/2	0.379 (x3x3)	15.4%
	Total		7.386	

Z DECAY MODES	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
$e^+ e^-$	(3.363 $\pm$ 0.004) %		45594
$\mu^+ \mu^-$	(3.366 $\pm$ 0.007) %		45594
$\tau^+ \tau^-$	(3.370 $\pm$ 0.008) %		45559
$\ell^+ \ell^-$	[b] (3.3658 $\pm$ 0.0023) %		—
$\ell^+ \ell^- \ell^+ \ell^-$	[h] (3.30 $\pm$ 0.31) $\times 10^{-6}$	S=1.1	45594
invisible	(20.00 $\pm$ 0.06) %		—
hadrons	(69.91 $\pm$ 0.06) %		—
$(u\bar{u} + c\bar{c})/2$	(11.6 $\pm$ 0.6) %		—
$(d\bar{d} + s\bar{s} + b\bar{b})/3$	(15.6 $\pm$ 0.4) %		—
$c\bar{c}$	(12.03 $\pm$ 0.21) %		—
$b\bar{b}$	(15.12 $\pm$ 0.05) %		—

$$\sin^2 \theta_W = 0.23126$$

don't forget the
factor of 3 for color
in quarks!

That's it!

- Please remember to respond to doodle poll for review session
 - I would like to set a time/place by the end of the week.
- Please send me any suggestions/topics/questions/issues for the review
- I will announce how you can collect PS 4