



Weekly Update

Shaghayegh Atashi July 10, 2017

Outline

- Some info on the partonic production channel
- How BdNMC calculates # of signal events
- Reproducing some plots from paper

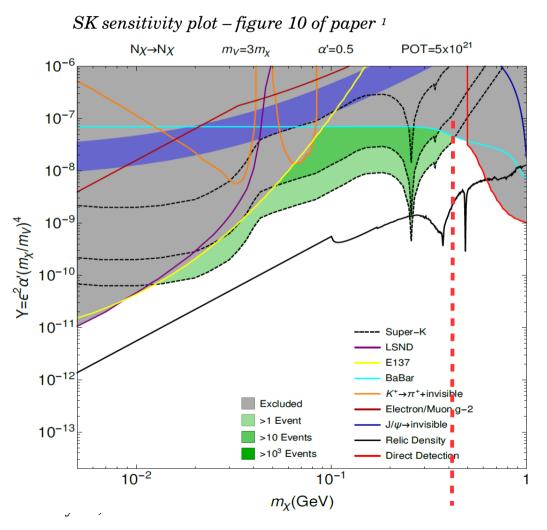
Update on parton_production channel

- * Reminder: corresponds to the dark matter production process $p+N \to V^* \to \chi/\chi^{\dagger}$
- Need 2 externally generated cvs files
 - Patrick can make these

Update on parton_production channel cont.

Question: Is this production channel relevant for SK?

• Partonic production relevant for $m_V > 1 GeV^{1}$



(with $m_V=3m_x$, alpha'=0.5, POT = $5x10^{21}$) sensitive up to $m_X \sim 0.7$ GeV

 $=> m_{\rm V} < 2.1~{\rm GeV}$

Couldn't find dependence of partonic production on beam energy ...

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How BdNMC calculates # of signal events

- I will show an overview of the simulation loop and how the # of signal events is calculated
- Before the simulation loop proper begins, a burn-in run is conducted for each production channel to estimate pmax, the maximum scattering probability encountered.
 - Very similar to the simulation loop, but the end state interaction results are not generated, and the variable nburn is incremented each time. The burn-in run is conducted until burntrials is equal to BURNMAX.

An overview of the simulation loop:

- samplesize is a parameter chosen by the user
- The code conducts trials until samplesize scattering events occur inside the detector
- ninteractions records the # of scattering events inside the detector, for each production channel

More detailed overview of the simulation loop:

 The total number of dark matter particles produced by each production channel is first calculated

1. Set trials=0, nevents=0, ninteractions = zeros(chan_count).

Initialize everything

2. While nevent < samplesize:

- (a) trials++.
- (b) Generate a uniform random number vrnd ∈ [0,vnumtot].
- (c) Set i = 0, scatterswitch=False.
- (d) While $i < \text{chan_count}$:
 - i. If vrnd < vnum[i] then break
 - ii. Else set vrnd = vrnd vnum[i].
- (e) Initialize Particle part with four-momentum p=0.
- (f) Set the four-momentum of part using Distribution's Sample_Particle method.
- (g) Initialize an empty list of Particles partlist.
- (h) Append part to partlist.
- (i) Generate a list of decay product Particles of part using DMGenerator_i's DMGen member function and store them in partlist.
- (j) For each dark matter Particle j in partlist:
 - i. If $Ldet_i == 0$ then continue
 - ii. Else
 - A. Simulate an interaction using Scatter's probscatter method.
 - B. If probscatter returns true, then insert the end state Particle generated by probscatter after j in partlist, and set scatterswitch=True.
- (k) If scatterswitch, then write all the particles in partlist to output_file, nevent++ and ninteraction $_i$ ++.
- 3. end while loop

- The total number of dark matter particles produced by each production channel is first calculated
- 1. Set trials=0, nevents=0, ninteractions = zeros(chan_count).
- 2. While nevent ≤ samplesize:
 - (a) trials++.
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Loop runs until number of scattering events (nevent) equal samplesize

- For each trial, nevent increments if at least one DM scatters inside the detector

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- 2. While nevent \leq samplesize:
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 - (b) Generate a uniform random number vrnd ∈ [0,vnumtot].
 - (c) Set i = 0, scatterswitch=False.
 - (d) While i < chan_count:</p>
 - i. If vrnd < vnum[i] then break
 - ii. Else set vrnd = vrnd vnum[i].

Choose a production channel

prob choosing ch. i = # DM particles produced by channel i

total # DM particles produced

- (e) Initialize Particle part with four-momentum p=0.
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The Distribution corresponding to that production channel generates the 4-momenta of a meson/V from the p-target collision via rejection sampling using a differential particle production cross-section or by iterating through a pre-generated list of particle 4-momenta

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 - (h) Append part to partlist.

Append the initial meson/V to partlist

- Generate a list of decay product Particles of part using DMGenerator, 's DMGen member function and store them in partlist.
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DMGenerator decays the initial V/meson, outputting a pair of dark matter particle 4-momenta.

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 - A. Simulate an interaction using Scatter's probscatter method.
 - B. If probscatter returns true, then insert the end state Particle generated by probscatter after j in partlist, and set scatterswitch=True.
 - (k) If scatterswitch, then write all the particles in partlist to output_file, nevent++ and ninteraction_i++
- 3. end while loop

If neither DM particle intersects the detector:

- choose another production channel

If either DM particle intersects the detector (the length of the intersection between the DM trajectories and the detector is nonzero):

- Determine whether the DM scatters.

If yes: an acceptance rejection algorithm is used to generate a final state 4-momentum by sampling from a differential distribution of interaction cross sections. Insert the end state Particle generated in partlist and set scatterswitch = True

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- 1. Set trials=0, nevents=0, ninteractions = zeros(chan_count).
- 2. While nevent < samplesize:
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 - (k) If scatterswitch, then write all the particles in partlist to output_file, nevent++ and ninteraction $_i$ ++.
- 3. end while loop

If neither DM particle intersects the detector:

- choose another production channel
- Does this mean we move on to the next trial and choose another production channel (i.e. trials is incremented)?
- If not (during a trial, we choose production channels until the DM intersects the detector), how do we account for the (small) probability of a DM intersecting the detector?

Eg: in one run of BdNMC (Run 1499450926):

vnum[i] = # DM particles produced by this channel = 4.32 e12

- **# DM intersecting detector = 1576**
 - → Will look at BdNMC code

If either DM particle intersects the detector (the length of the intersection between the DM trajectories and the detector is nonzero):

- Determine whether the DM scatters.

If yes: an acceptance rejection algorithm is used to generate a final state 4-momentum by sampling from a differential distribution of interaction cross sections. Insert the end state Particle generated in partlist and set scatterswitch = True 14

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 - (k) If scatterswitch, then write all the particles in partlist to output_file, nevent++ and ninteraction $_i$ ++.
- 3. end while loop

If either particle scatters, then write all particles in partlist to the output_file, increment nevent, and increment ninteractions for the appropriate production channel

If neither particle scattered, then move on to the next trial

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 - (k) If scatterswitch, then write all the particles in partlist to output_file, nevent++ and ninteraction_i++.

3. end while loop

End the while loop when nevent = samplesize

At this point, the number of scattering from each production channel i is stored in ninteraction_i

How BdNMC calculates # of signal events

For each production channel i:

$$\texttt{signal_events}[i] = \frac{\texttt{ninteractions}[i]}{\texttt{trials}} \times \texttt{vnumtot} \times \texttt{pmax} \times \texttt{efficiency},$$

Note: x trials are done until samplesize total scattering events occur, where samplesize if an input parameter

- Sum of ninteractions[i] over all production channels = samplesize
- ninteractions[i] = number of scatterings by DM originating from production channel i
 vnumtot is the total number of DM particles produced
- Pmax is "the maximum scattering probaility"
- Efficiency = detector efficiency
- Checked this formula by hand using the values for trials, vnumtot, etc outputed by BdNMC in terminal
- Why do we multiply by pmax? Isn't the scattering probability taken care of in the trials?

Reproducing figure 1.0 of "Light dark matter in neutrino beams: production modelling and scattering signatures at MiniBooNE, T2K and SHiP":

- I run BdNMC ~20 times for m_V in [0.005, 0.75] for the production channels pi0 decay, eta decay, and proton bremsstrahlung (can't easily do parton production, but it doesn't matter for m_v<~0.75GeV) and signal channel NCE_nucleon.
- Parameters:

$$m_{\chi} = 0.01 \text{ GeV}, \ \epsilon = 10^{-3} \text{ and } \alpha' = 0.1.$$

- Figure's caption says plot is for 9 GeV beam energy, but it's actually 8.9 GeV, they rounded.
- 2e20 POT
- Efficiency = 0.35
- 0.9 pi0_per_POT (pi0_per_POT is the number of pi0's expected per proton on target)
- Using production distributions pi0_sanfordwang for pi0 decay and k0_sanfordwang for eta decay (Sanford-Wang distributions appropriate for MiniBooNE energies)
- * Production distribution proton_brem and ptmax = 0.2, zmin = 0.3, zmax = 0.7, which are appropriate for MiniBooNE
- Continued on next slide...

Reproducing some plots from the paper

• Want to reproduce figure 1.0 of "Light dark matter in neutrino beams: production modelling and scattering signatures at MiniBooNE, T2K and SHiP":

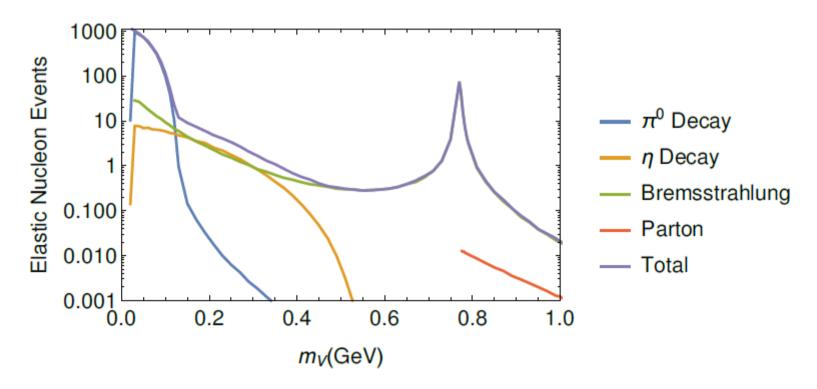


FIG. 1. A plot illustrating the distinct contributions to DM production (coupled through the vector portal), as discussed in the text, using the 9 GeV proton beam at MiniBooNE as an example. The rate of elastic scattering events on nucleons is plotted versus the vector mediator mass. From smaller to larger values of m_V , the dominant channels are π^0 decays, η decay, bremsstrahlung, which becomes resonant near the ρ/ω mass region, and finally direct parton-level production. The plot uses $m_{\chi} = 0.01 \text{ GeV}$, $\epsilon = 10^{-3}$ and $\alpha' = 0.1$.

Reproducing figure 1.0 of "Light dark matter in neutrino beams: production modelling and scattering signatures at MiniBooNE, T2K and SHiP":

- Cuts on the kinetic energy of outgoing nucleon, default to min=0 and max=1e9 GeV:
- Initially Patrick said he used min=0.35 GeV and max=1 GeV
- Sensitivity section of paper says they use the same cuts (page 10)
- Using these cuts, my plot was off (~100 times fewer events than figure 1.0)
- I saw the paper "Dark Matter Search in a Proton Beam Dump with MiniBooNE" (arXiv:1702.02688v2) used 0.035 GeV as the lower cut in their analysis, do I used this as the lower cut
 - No longer off by a factor of 100, still a bit off.
 - Asked Patrick → 0.35 is a type, **it is actually 0.035GeV!** (upper energy cut is still 1.0 GeV)
- I'm also using the same value for the detector geometry as Patrick used (a bit different than values in the tables in paper):

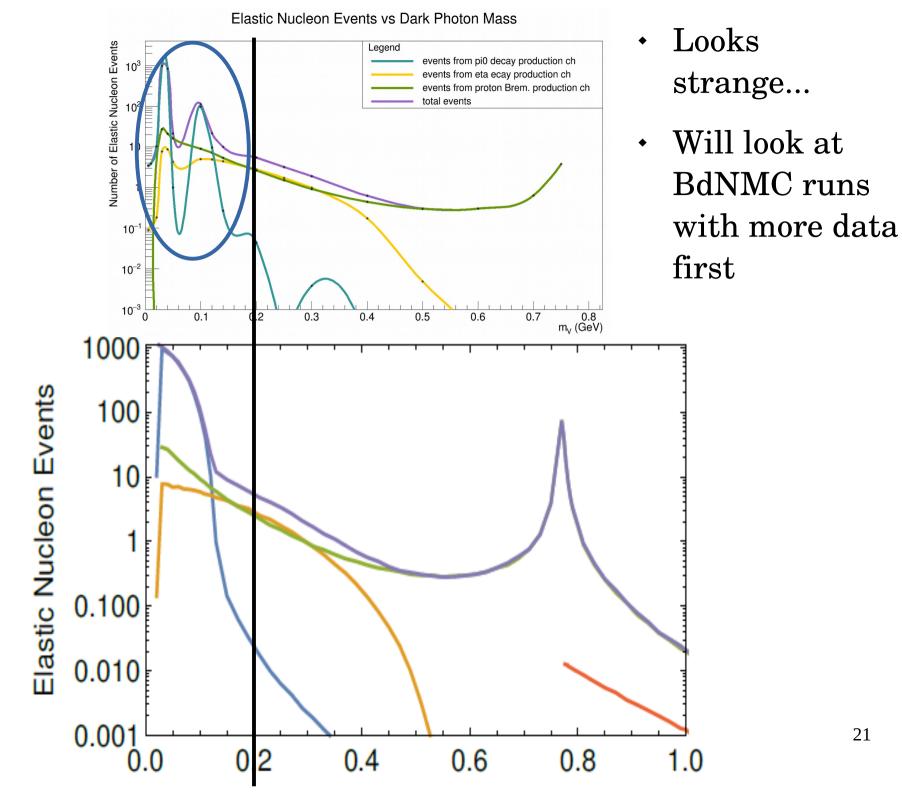
$$x = 0.0 \text{ m}$$

$$y = -1.9 \text{ m}$$

$$z = 491 \text{ m}$$

$$r = 5 m$$

• My plot is on the next page (only 1000 entries for each BdNMC run, am currently running it with 20000 entries but that will take a few hours)



Backup slides

A bit more info on the proton Bremsstrahlung production channel:

- User needs to state the values of these parameters when using the proton brem. channel:

```
#This invokes the bremsstrahlung production channel. This works, but may be #unreliable around the rho resonance. The zmin/zmax values seem reasonable #for MiniBooNE energies. ptmax could be as large as the proton mass, but #probably would not change signal much. production_channel V_decay production_distribution proton_brem ptmax 0.2 zmin 0 3 zmax 1.7
```

ptmax: The maximum transverse momentum which a produced V mediator may possess. The minimum is assumed to be 0.

zmin: The minimum value of $z = \frac{p_{V,z}}{P}$, where $p_{V,z}$ is the momentum of the V parallel to the z axis, and P is the total momentum of a beam proton incident on the target.

zmax: The maximum value of z, defined as in the **zmin**.

Calculating the total number of V bosons (N_v) produced by the proton bremsstrahlung production channel: (V mediators decay to dark matter isotropically)

$$N_V = \mathbf{POT} \int_0^{\mathbf{ptmax}^2} dp_\perp^2 \int_{\mathbf{zmin}}^{\mathbf{zmax}} dz \frac{d^2 N_V}{dz dp_\perp^2}, \tag{11}$$

Where the differential V production rate is:

$$\frac{d^2 N_V}{dz dp_\perp^2} = \frac{\sigma_{pA}(s')}{\sigma_{pA}(s)} F_{1,N}^2(q^2) w_{ba}(z, p_\perp^2), \tag{6}$$

where $s' = 2m_p(E_p - E_V)$, $s = 2m_pE_p$ and the photon splitting function is [54]

$$w_{ba}(z, p_{\perp}^{2}) = \frac{k_{V,B}^{(0)}}{2\pi H} \left[\frac{1 + (1-z)^{2}}{z} - 2z(1-z) \left(\frac{2m_{p}^{2} + m_{V}^{2}}{H} - z^{2} \frac{2m_{p}^{4}}{H^{2}} \right) + 2z(1-z)(z + (1-z)^{2}) \frac{m_{p}^{2} m_{V}^{2}}{H^{2}} + 2z(1-z)^{2} \frac{m_{V}^{4}}{H^{2}} \right],$$

with $H = p_{\perp}^2 + (1-z)m_V^2 + z^2m_p^2$, and $k_{V,B}^{(n)}$ was defined above in (5). $\longrightarrow k_{V,B}^{(n)} = \begin{cases} \epsilon^2\alpha(\alpha')^n & \text{for } U(1)' \\ \alpha_B^{n+1} & \text{for } U(1)_B \end{cases}$

Proton Bremsstrahlung Production Channel cont.

- . The paper goes into detail about the "timelike form factor $F_{1,N}(q^2)$ (page 5 of "Light dark matter in neutrino beams: production modelling and scattering signatures at MiniBooNE, T2K and ShiP")
- → not really sure what they mean, but basically: the equation that the paper uses for the differential V production rate requires the following kinematic conditions:

$$E_P >> m_P$$

 $E_V >> m_V$
 $E_P - E_V >> |P_L|$

- So the range of values of z and $|P_{\perp}|$ over which the integral for N_{V} is calculated must be such that the conditions are met.
- Paper says these conditions are met if $z \in [0.1, 0.9]$ and $|P_{\perp}| < 1$ GeV for high energy experiments like SHiP, but not for lower energy experiments.
- eg for MiniBooNE (much lower energy than SHiP), the paper uses $z \in [0.3, 0.7]$ and $|P_1| < 0.2$ GeV (as in the parameter cards for the MiniBooNE-like experiments)

